

Image Contrast Optimization of the CE-FAST Pulse Sequence on a 0.05T Imager

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Received October 5, 1995

In this work we present experiments carried out on our 0.05T imager in order to optimize the image contrast of the CE-Fast pulse sequence. We calculated the expected contrast for the CE-FAST and Spin Echo sequences and compare these with values obtained from images of a phantom built for this purpose. We conclude that the CE-FAST sequence is superior to the traditional Spin Echo technique if one is interested in producing a single image with T_2 contrast for the fast survey of a region.

I. Introduction

Different rapid pulse sequences that take advantage of the nuclear magnetization steady-state condition^[1] have been proposed for fast acquisition of MRI images. Some of these sequences have already been investigated in detail^[2-6] and are now being used routinely for clinical purposes such as heart-triggered imaging as well as for other applications like functional imaging and localization images for *in vivo* spectroscopy.

We have implemented in our home made 0.05T whole body imaging system, a pulse sequence that combines two steady-state techniques. This pulse sequence, originally proposed by different authors^[7-9], permits the simultaneous acquisition of two images with clearly different contrasts. The signal from the first acquisition interval, the FID, is the same as for the standard Fourier Acquired Steady State (FAST) sequence, which is also known as Fast Field Echo (FFE). The second part of the signal, the echo, which is located immediately before the radio frequency (RF) pulses, is analogous to the signal presented by the sequence known as contrast enhanced FAST (CE-FAST) or T_2 weighted FFE^[2-5]. In this work we will focus on the technique that uses this second signal (CE-FAST). Images obtained using this technique can show high T_2 contrast when properly optimized. In the following we will dis-

cuss the contrast optimization of this sequence.

Using steady-state techniques the contrast of the images is basically controlled by varying the repetition time (TR) and the flip angle (α). The echo time (TE), which is used conventionally for contrast control, is generally kept short in steady-state techniques to minimize inhomogeneity and susceptibility effects that arise due to the use of gradient recalled echoes. The contrast of the images obtained with steady-state techniques is very sensitive to TR and α ^[5,6]. In particular, for appropriate TR values the CE-FAST sequence presents intrinsically high T_2 contrast. This may be understood in a simplified model regarding the acquired echo as the result of two consecutive RF pulses with an effective echo time of $TE_{\text{eff}} = 2 \cdot TR - TE$, where TE is the interval between the echo and the next RF pulse.

II. Fundamental equations

The pulse sequence for CE-FAST is given in Fig. 1 where the relevant sequence parameters are defined.

In order to analyze theoretically the contrast of the CE-FAST pulse sequence we have to know the analytical expression for the signal amplitude, which is proportional to the transverse magnetization M_Y . The expression of M_Y in the steady-state condition in the absence of imaging gradients has been given by Ernst et. al.^[1] as

$$M_Y = \frac{M_0(1 - E_1)E_2 \sin \alpha (\cos O - E_2)}{(1 - E_1 \cos \alpha)(1 - E_2 \cos O) - (E_1 - \cos \alpha)(E_2 - \cos O)E_2} \quad (1a)$$

with

$$E_1 = e^{-TR/T_1} \quad (1b)$$

and

$$E_2 = e^{-TR/T_2}, \quad (1c)$$

where O is the total precession angle of the spins during one repetition cycle. In order to avoid artifacts, the gradients usually applied for imaging are large enough to produce a rapid spatial variation of O so that, within a single voxel, O must be replaced by its mean value^[4,10]. This means that for calculations Eq. (1a) has to be integrated over O . The analytical result of this integration was given by van der Meulen et al.^[2] as

$$M_Y = \frac{M_0 \sin \alpha}{1 + \cos \alpha} \left[\frac{(1 + \cos \alpha)E_2^2 - a}{\sqrt{a^2 - b^2}} + 1 \right] e^{(TR-TE)/T_2} \quad (2a)$$

with

$$a = \frac{1 - E_1 E_2^2 + \cos \alpha (E_2^2 - E_1)}{1 - E_1}, \quad (2b)$$

$$b = (1 + \cos \alpha)E_2. \quad (2c)$$

The imaging gradients as shown in figure 1 shift the echo signal from TR to TE (for this reason the factor $e^{(TR-TE)/T_2}$ has been included in Eq.(2a) as in Ref.[2]).

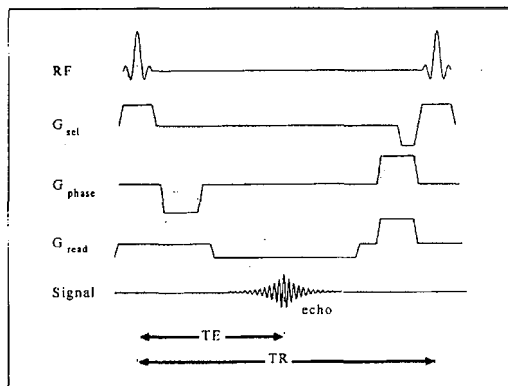


Figure 1. CE-FAST pulse sequence.

We now may calculate the contrast between tissues with different relaxation times T_1 , T_2 and proton

spin densities which are proportional to the equilibrium magnetization M_0 . For long repetition times the sequence gives intrinsically high T_2 contrast, though the signal amplitude and therefore signal to noise ratio decays in an exponential manner. The same applies to the flip angle where the contrast is in general higher for small angles. A reasonable compromise between contrast and signal to noise is needed. The definition of a contrast-to-noise ratio (CNR) is very useful for this purpose^[11,12].

For two tissues A and B with their respective signal amplitudes S_A , S_B and standard deviations σ_A , σ_B we calculate the CNR as^[13]

$$CNR = \frac{S_A - S_B}{k}, \quad (3a)$$

with

$$k = \sqrt{(\sigma_A^2 + \sigma_B^2)/2}. \quad (3b)$$

III. Contrast calculations

Using Eq. 2, we calculated the theoretical contrast for the CE-FAST sequence for two hypothetical tissues, A and B, having M_0 equal to unity and relaxation times T_1 and T_2 as given below. Fig. 2 shows the results as contour plots of $S_A - S_B$ versus repetition time, TR , and flip angle, α . From this one can see that an absolute maximum contrast can be achieved, between two given tissues, by carefully optimizing the sequence parameters for typical values of the corresponding relaxation times.

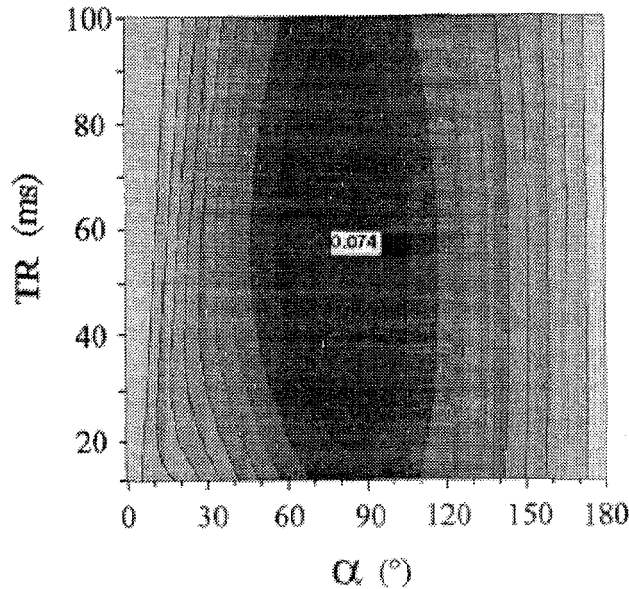


Figure 2. Contour map of the contrast between tissues A and B for the CE-FAST sequence. The maximum is at $TR = 56\text{ms}$, $\alpha = 79^\circ$ with a value of 0.074.

In order to prove that the above procedure allows to correctly optimize the sequence we performed a series of experiments using a phantom built for this purpose. The phantom consists of two concentric tubes: the outer one contains CuSO_4 solution with a small amount of MnCl_2 , the concentrations being 0.5mM and 0.05mM respectively, and the inner tube is filled with a 0.5mM MnCl_2 . These solutions, A and B, had the same relaxation times respectively as those assumed above for the two hypothetical tissues A and B, given in Table 1. In the experiments we varied the flip angle from 24° to 136° for four different repetition times. The results are shown in Fig. 3 as the difference between the mean values over regions of interest of the same area for solutions A and B of the non normalized images. The error bars represent the quantity k in Eq. 3b obtained from the standard deviations of the intensities over the same regions. The continuous lines in Fig. 3 are plots of the expected contrast as calculated from Eq. 2, where the value of M_0 was obtained from a series of images using the spin echo sequence with long TR and varying TE. This was done by extrapolating the fitted image intensities to $TE = 0$. In this way the gain of the instrument is taken care of and the figure allows to check the accuracy of Eq. 2.

In all the above experiments the flip angles were measured using a direct method based on a stimulated echo experiment as described in a previous paper^[14].

Table 1: Relaxation times of tissues A and B.

T_{1A} (ms)	T_{2A} (ms)	T_{1B} (ms)	T_{2B} (ms)
273 ± 14	162 ± 1	237 ± 9	86 ± 1

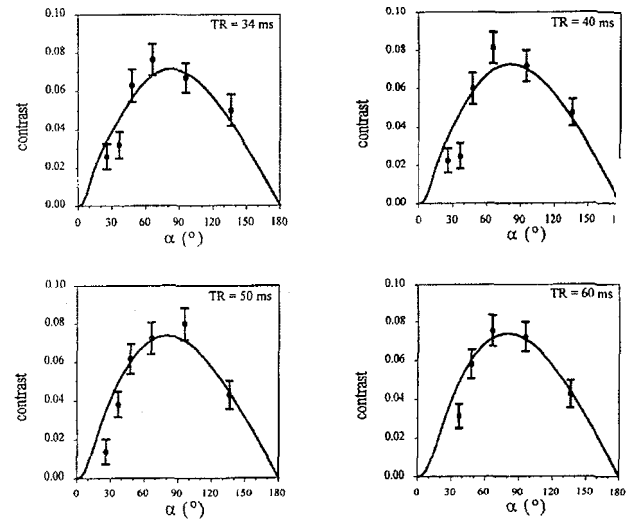


Figure 3. Comparison between experimental and calculated contrast for the CE-FAST sequence using different repetition times. The asterisks show the measured difference between the mean values over regions of interest of the same area for solutions A and B of the non normalized images. The error bars represent the quantity k in Eq. 3b obtained from the standard deviations of the intensities over the same regions. The continuous lines are plots of the expected contrast as calculated from Eq. 2.

IV. Comparison between CE-FAST and Spin Echo

The above calculations give the optimum TR for maximum contrast, $S_A - S_B$. In order to optimize the contrast to noise ratio that can be obtained in a given time it is necessary to take into account the fact that reducing TR allows to increase the number of averages and therefore to reduce the noise. To account for this we have to divide the calculated value of $S_A - S_B$ by the square root of TR. We have done this for both, the CE-FAST and the Spin Echo sequences, and the results are plotted in Fig. 4 below. The figure shows that for the shortest TR that could be realized on our system, 34 ms, it is possible to expect better contrast with a CE-FAST sequence than with optimized spin echo. This is

particularly significant if one is interested in producing a single image with T_2 contrast for the fast survey of a region.

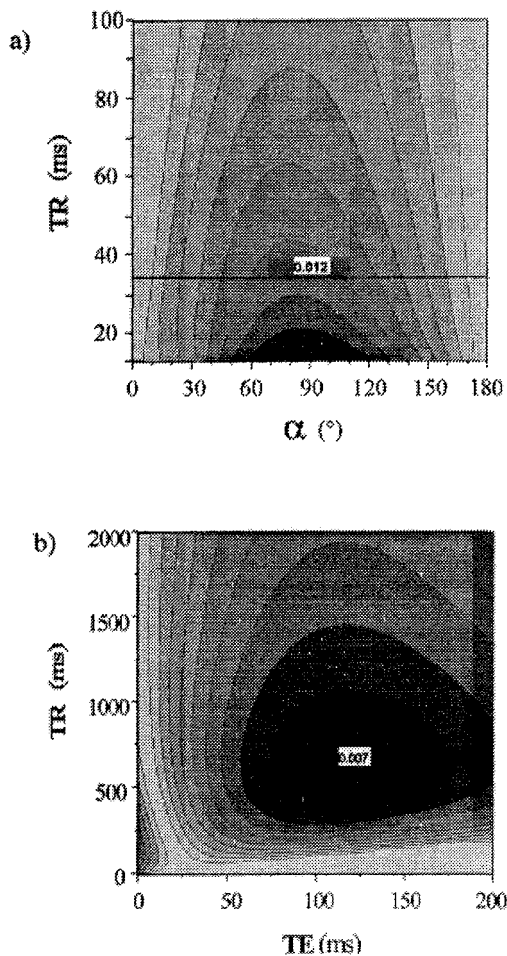


Figure 4. Contour plots of the expected contrast per unit time: a) for the CE-FAST sequence the maximum occurs for $TR=34\text{ms}$, $\alpha = 82^\circ$ and equals 0.012; b) for the Spin Echo sequence, the maximum is at $TR=617\text{ms}$, $TE = 116\text{ms}$ with a value of 0.007.

In order to see how the two sequences compare experimentally in our 0.05 Tesla scanner we have imaged our previously described phantom with both methods. The resulting images are given in Fig. 5, showing the imaging parameters used, the resulting total acquisition times and the Contrast to Noise ratio, C/N , measured in the images using the definition given in Eq. 3. One optimized CE-FAST and two different Spin Echo images, all acquired with the same total imaging time $T_{tot}=2':45''$, show the superior result obtained with the first when short acquisition of a single slice is desired. We have so included, as a reference, the image obtained

with our standard T2 Spin Echo multislice protocol, which uses 4 averages and takes 17':00" minutes giving practically the same C/N ratio, as the CE-FAST image.

We want to point out that the above discussion was based on the assumption that in a given acquisition time only one image is being acquired, which is one restriction of the CE-FAST sequence. However with the spin echo technique it is possible to acquire multiple images of slices in different positions of the sample using interleaved acquisition. When using CE-FAST for multiple slices the experiment normally has to be executed separately for each image, which increases the total acquisition time significantly. Nevertheless our results are valid for the case that only one image is needed and maximum CNR is of concern as for example in a first scan for localization of a subsequent complete MRI or MRS experiment.

V. Conclusion

From our measurements on a doped water phantom we proved that CNR calculations based on the relaxation times of two tissues are in agreement with experimentally determined CNR values. It is therefore possible to find the maximum CNR of two tissues by calculation solely knowing their relaxation times and proton spin densities which, in many cases, are readily available in literature. When tissue parameters are not available the CNR can easily be optimized choosing the lowest possible repetition time and maximizing the CNR with respect to the flip angle α .

We conclude that the CE-FAST sequence, with adequate experimental settings and even at a low field of 0.05T, is a possible alternative to the standard spin echo technique when highly contrasted single scans are needed, such as for example in a first scan for localization of a subsequent complete MRI or MRS experiment.

Acknowledgment

We would like to thank Eng. Edson L. G. Vidoto, Eng. M. J. Martins and Mr. B. Muniz for their help with experiments and discussion.

This work was supported in part by the Brazilian institutes for scientific research FAPESP, FINEP, PADCT and CNPq.

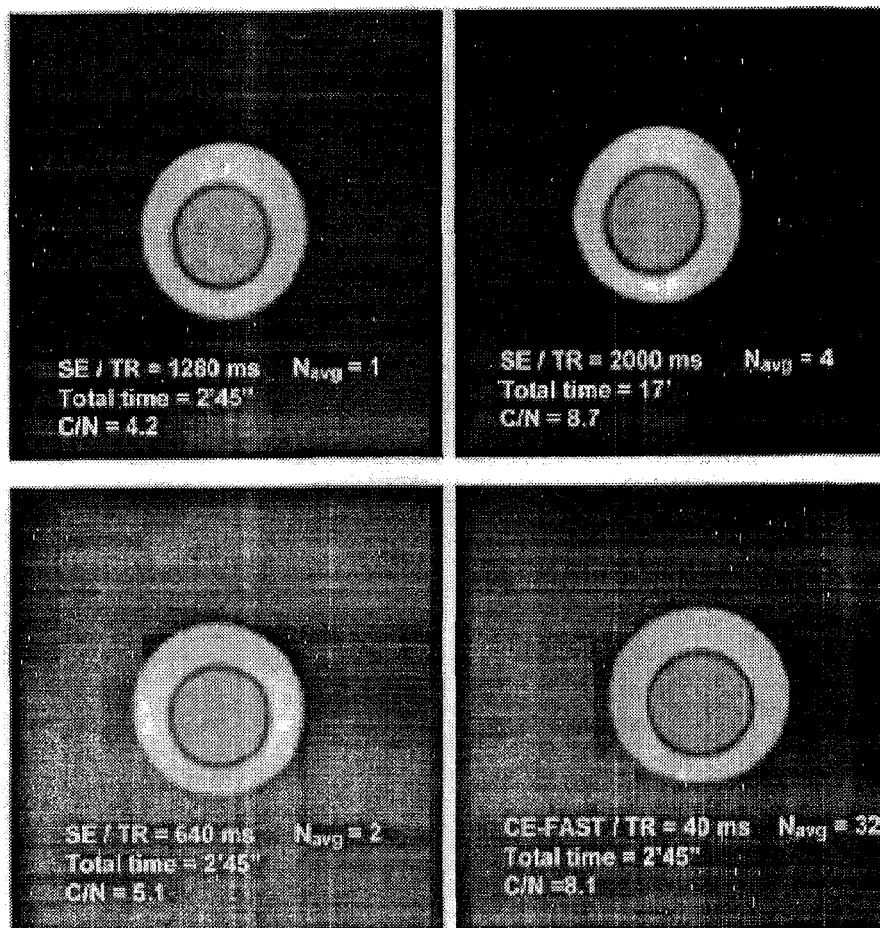


Figure 5. Comparison between CE-FAST and Spin Echo techniques. The two bottom images show that CE-FAST is superior for quick single slice acquisition. The two images on the right show that CE-FAST can achieve the same C/N ratio as the standard T2 weighted Spin Echo, in a much shorter time. The two images on the left show the improvement in C/N, for equal acquisition times, obtained by choosing a compatible TR close to the optimum value according to Fig. 4.

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