

Hydrogen Motion and Light-Induced Defects in Hydrogenated Amorphous Silicon

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Hydrogen diffusion experiments in hydrogenated amorphous silicon (a-Si:H) demonstrate that hydrogen migration is enhanced under illumination. The enhancement is suppressed when the photo-generated carriers are extracted from the diffusion region by an electric field. The dependence of hydrogen diffusion on illumination suggests a close connection between hydrogen diffusion and metastability in a-Si:H.

I. Introduction

Hydrogenated amorphous silicon (a-Si:H) is a promising material for application in large area electronics and in photovoltaic energy conversion. a-Si:H films of good electronic properties can be deposited over large areas by standard thin-film deposition processes such as sputtering of a silicon target and plasma enhanced decomposition of silane gas (SiH_4). The material has a large light absorption coefficient in the visible range and has been used as active layers in field-effect transistors, photo-detectors, and solar cells.

Unfortunately, the electronic properties of a-Si:H degrade when extra electronic carriers (i.e., electrons and/or holes) are introduced in the material by illumination^[1], current^[2] or the application of an electric field^[3]. The degradation is reversible and can be eliminated by annealing the material above $\sim 180^\circ\text{C}$. The degradation is associated with the creation of defects with electronic states near the center of the energy gap. The defects have been identified as unsaturated silicon bonds (dangling-bonds) which act as effective recombination centers for minority carriers and considerably decrease the doping efficiency of the material. A considerable effort has been devoted in the recent years to obtain films of high stability against defect formation^[4].

Two classes of models have been proposed to explain the origin of metastable dangling-bonds in a-Si:H. In the first class ("extrinsic models") dangling-bond formation is attributed to the weakening of silicon bonds

in the neighborhood of an extrinsic impurity atom^[5]. Many impurity atoms (such as, for instance, oxygen and carbon) are present in a-Si:H films in concentrations well exceeding the density of metastable defects ($\sim 10^{17} \text{ cm}^{-3}$). In addition, it is well-established that the incorporation of dopant atoms in a-Si:H is accompanied by an increase in the defect density^[6].

In the second class of models ("intrinsic models"), on the other hand, metastability is viewed as an intrinsic property of the a-Si:H network. It has been proposed that the presence of excess carriers weakens silicon-hydrogen^[7] or neighboring silicon-silicon bonds^[8] and increases the emission probability of a hydrogen atom from a silicon-hydrogen bond, leading to the formation of a dangling bond. If the emitted atom migrates away, a silicon dangling-bond remains when the carrier population is reduced to the thermal equilibrium value after the removal of the external excitation^[9]. The metastable character of dangling-bonds arise from the small hydrogen migration rate around room temperature. At higher temperatures (i.e., above $\sim 180^\circ\text{C}$) hydrogen diffuses relatively fast through the network so that the defect can be quickly annihilated. In this regime the amorphous network responds quickly to changes in the carrier concentration and there is a thermal equilibrium between carriers and defects.

A natural consequence of the intrinsic model for metastability is that hydrogen diffusion should be enhanced when extra carriers are introduced in a-Si:H. In this paper we review the experimental results on hydrogen diffusion and its dependence on carrier concentration in a-Si:H. As will be described in Sec. II, the ex-

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periments were performed by diffusing deuterium from a plasma or from a deuterated $a\text{-Si:H:D}$ layer into an intrinsic $a\text{-Si:H}$ film. The deuterium profiles after diffusion were determined by Secondary Ion Mass Spectrometry (SIMS). In Sec. III we present experimental results demonstrating that hydrogen migration in $a\text{-Si:H}$ is controlled by an electronic mechanism and (i) is enhanced when the carrier population is increased by illumination and (ii) is suppressed when the photo-generated carriers are extracted from the diffusion region by an electric field. Microscopic models for hydrogen migration and its connection to metastability are discussed in Sec. IV.

II. Experimental

Light-enhanced hydrogen diffusion was investigated by diffusing deuterium (^2H) into an $a\text{-Si:H}$ film during illumination and measuring the resulting deuterium profile by secondary ion mass spectrometry (SIMS)^[10]. The deuterium diffusion source was either a plasma or a deuterated amorphous silicon layer ($a\text{-Si:H:D}$). In the first case (see Fig. 1a), single glow-discharge $a\text{-Si:H}$ layers were exposed to a remote microwave plasma in the dark or under illumination from an Ar-laser ($\sim 15 \text{ W/cm}^2$). In the second case, the experiments were performed in the $a\text{-Si:H}/a\text{-Si:H:D}/a\text{-Si:H}$ multilayer structures shown in Fig. 1b. These structures have layer thicknesses of 100nm, 200nm, and 500 nm, respectively. The bottom $a\text{-Si:H}$ film was grown by glow discharge decomposition of pure silane at 230°C . The two capping layers were deposited at a lower temperature (130°C) in order to reduce deuterium diffusion during growth, and deuterium incorporation in the $a\text{-Si:H}$ layer was achieved by diluting silane with a 20% volume fraction of deuterium. The samples in this case were illuminated with water-filtered light from a xenon arc lamp (15 to 20 W/cm^2).

Diffusion experiments were also performed in p-i-n $a\text{-Si:H}$ photodiodes (Fig.1c). The deuterium source in this case is an $a\text{-Si:H:D}$ layer inserted in the middle of the intrinsic layer. All layers were deposited at 200°C . The diodes have transparent indium-tin oxide top contact to allow illumination of the intrinsic layer during the thermal treatment. The main advantage of using these photodiodes in diffusion studies is that the carrier concentration in the diffusion region can be controlled either by illumination or by an external bias^[11].

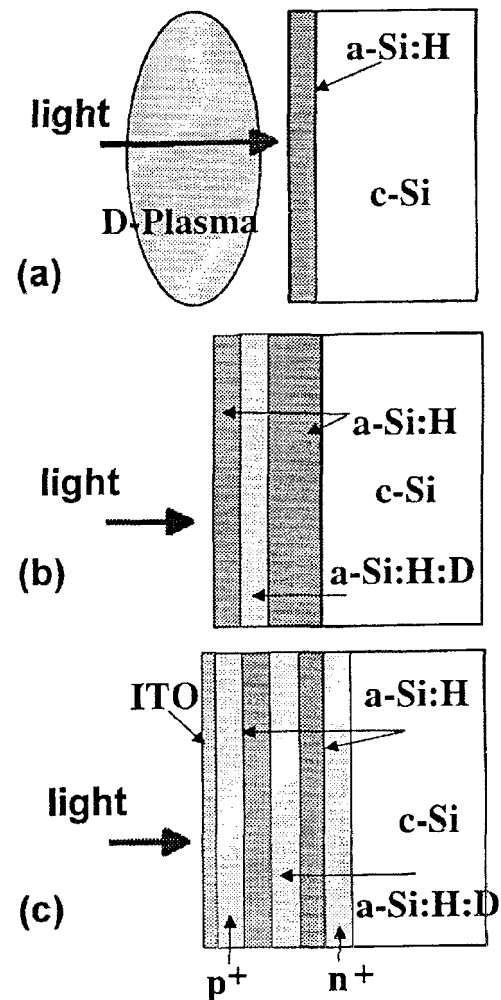


Figure 1: Deuterium diffusion from (a) a deuterium plasma and (b) from a deuterated $a\text{-Si:H}$. In (c) the deuterated layer was inserted into the intrinsic region of a p-i-n solar cell. All $a\text{-Si:H}$ layers were deposited by the glow discharge decomposition of silane. The solar cell is covered with a transparent indium-tin-oxide (ITO) film and a chromium grid (not shown in the diagram). In all cases the samples were grown on crystalline silicon substrates (c-Si).

In all experiments the samples were mounted on a temperature-controlled heating stage and special care was taken to reduce sample heating due to the strong light intensity ($15 - 20 \text{ W/cm}^2$) on the sample surface. The samples were grown on crystalline silicon substrate, and a heat conductive paste was used to minimize temperature gradients between the substrates and the temperature-controlled sample holders. By using the thermal paste the temperature increase (for samples measured in vacuum) due to light absorption was reduced from $\sim 20^\circ\text{C}$ (without paste) to less than 5°C (i.e., with paste the thermal component was negligible). As an additional precaution, the diffusion experiments

in multilayer structures were always performed by irradiating simultaneously two samples: one facing the lamp and the other ("dark") illuminated from the back side of the silicon substrate. The temperature difference between the illuminated and the non-illuminated samples is thereby reduced to the thermal gradient across the silicon wafer ($< 3^\circ\text{C}$).

III. Results

Light-enhanced deuterium diffusion^[12] is illustrated in Figs. 2a to 2c, where deuterium concentration profiles in multilayer structures annealed in the dark (thin lines) and under illumination (thick lines) at different temperatures are compared. The diffusion times are indicated in each case. The vertical line in the figure marks the interface between the a-Si:H:D and the bottom a-Si:H layers. The deuterium concentration drops from $7 \times 10^{19} \text{ cm}^{-3}$ in the a-Si:H:D layer, corresponding to an atomic concentration of 1.5% relative to the silicon atoms, to $6 \times 10^{17} \text{ cm}^{-3}$ in the a-Si:H layer. The latter concentration corresponds to the natural deuterium abundance in a-Si:H with 8% hydrogen. At 200°C and below (Fig. 2a), diffusion in the dark is negligible in comparison with the depth resolution of the SIMS profiles of $\sim 8.0 \text{ nm}$ for each factor of e decay in concentration. Under illumination, diffusion is strongly enhanced and can be noticed in the SIMS profiles for temperatures above $\sim 175^\circ\text{C}$. Diffusion in the dark and under illumination are temperature-activated, and the largest difference between the two profiles occur, at 250°C , for the illumination intensity used here.

The deuterium concentration profiles in the a-Si:H layers in Fig. 2a to 2c are well described by a complementary error function^[10], $\text{erfc}[x/2(D_{\text{eff}}t_d)^{1/2}]$. Here, x is the depth in the layer, and D_{eff} and t_d are, respectively, the effective diffusion coefficient and the diffusion time. The effective diffusion coefficient in the dark is thermally activated with an activation energy of 1.2-1.5 eV. Under illumination, the diffusion coefficient at 250°C increases by an order of magnitude, and the enhancement is even larger for lower temperatures.

Deuterium incorporation from the plasma in the dark (thin lines) and under illumination (thick lines) is illustrated in Fig. 2d. Enhanced deuterium diffusion is also observed in this case, but the diffusion coefficient increases by a factor of only 1.5 to 2, as compared to an order of magnitude enhancement observed in diffusion from a deuterated layer. In diffusion experiments from the plasma the atoms are directly injected from the gas

phase into weakly bonded interstitial sites. In contrast, when the deuterium atoms proceed from an a-Si:H:D layer, they first have to be emitted from Si-D bonds to the interstitial migration path. The diffusion kinetics is then determined by the deuterium emission rate from Si-D bonds^[13]. The larger diffusion enhancement in the latter case indicates that light-enhancement is associated with an increase in the deuterium emission rate in the presence of photo-generated carriers^[12,14].

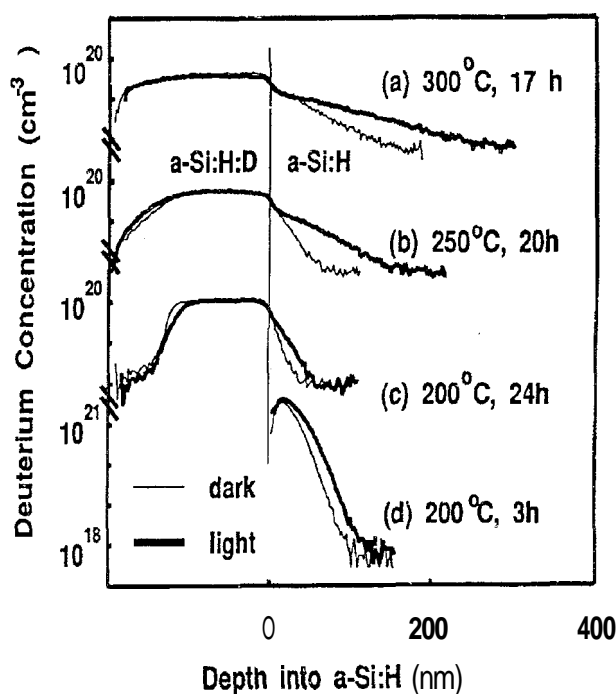


Figure 2: Deuterium depth profiles after diffusion in the dark (thin lines) and under illumination with 17 W/cm^2 from a xenon lamp (thick lines) in (a)-(c) a-Si:H/a-Si:H:D/a-Si:H multilayer. Curve (d) shows corresponding results for a-Si:H layers exposed to a remote deuterium plasma. In this case the samples were illuminated with the $\lambda = 514.5 \text{ nm}$ line of an argon laser. The diffusion time and temperature are indicated in each case. The vertical line marks the beginning of the bottom a-Si:H layers, where deuterium diffusion was studied.

Hydrogen diffusion under illumination was also investigated in doped and in compensated amorphous silicon films. In this study, the bottom a-Si:H layer in the multilayer structures (see Fig. 1 b) was replaced by a doped film. Figure 3 compares deuterium concentration profiles after annealing at 250°C in the dark (dashed lines) and under illumination (thick lines) for boron doped, phosphorus doped, and for a compensated film. The corresponding diffusion profiles for an undoped sample at the same temperature are shown

in Fig. 2b. The phosphorus (boron) doped layer was grown with a phosphine (diborane) to silane gas phase molar ratio of 10^{-4} (10^{-2}). For the compensated sample, equal molar ratios of 10^{-3} of phosphine and diborane were used. When compared to undoped material, the diffusion coefficient in the dark is approximately a factor of $\sim 5 - 10$ larger in the compensated and in the phosphorus doped samples, and a factor of 100 larger in the boron doped film. Under illumination, the diffusion coefficient increases by less than 20%, as compared to an order of magnitude increase in undoped samples. Since any temperature increase due to light absorption is expected to be the same in doped and undoped samples, these results clearly demonstrate that the large light-induced enhancement of the hydrogen diffusion coefficient in undoped a-Si:H is not due to sample heating during illumination.

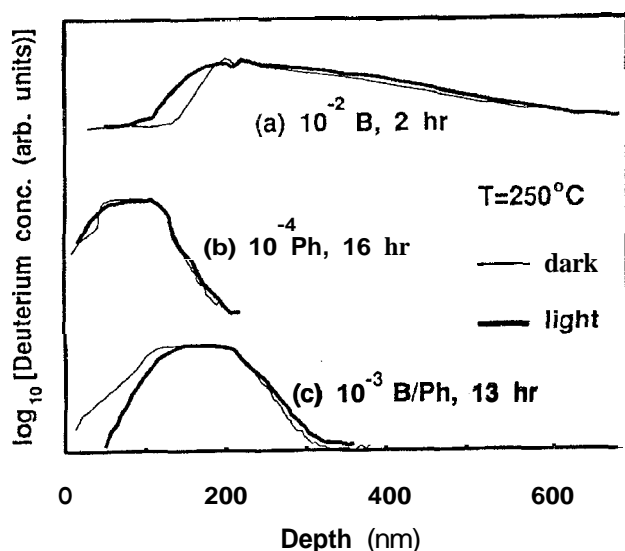


Figure 3: Diffusion profiles in the dark (thin lines) and under illumination with 17 W/cm^2 from a xenon lamp (thick lines) in (a) boron-doped (1% volume ratio of diborane in silane), (b) phosphorus-doped (0.01% volume ratio of phosphine in silane), and in (c) compensated a-Si:H layers (0.1% phosphine plus 0.1% diborane in silane). The diffusion experiments were performed at 250°C for the times indicated in the figure.

The interaction between hydrogen migration and electronic carriers was further investigated by performing diffusion experiments where the a-Si:H/a-Si:H:D/a-Si:H multilayers were inserted in the intrinsic region (i-layer) of a p-i-n solar cell (see Fig. 1c). The carrier concentration in the diffusion region in this case can be controlled by varying either the external bias or the

illumination intensity. An example is shown in Fig. 4. The dashed line in this figure is the deuterium profile in a control sample annealed at 225°C in the dark. The thin solid line reproduces the profile obtained when the solar cell is illuminated during the thermal treatment. As in Fig. 2, the profiles are broadened at the interfaces due to light-induced diffusion. The thick line in Fig. 4 shows the combined effects of illumination and of an average reverse field of $9.2 \times 10^4 \text{ Wcm}$ in the intrinsic layer during annealing obtained by applying a reverse bias of 5V. The diffusion profiles are asymmetric with respect to the center of the a-Si:H:D layer due to an inhomogeneous carrier density in the intrinsic layer, and small changes are observed on the right-hand side of the deuterated layer^[15]. On the left-hand side, on the other hand, light-enhancement of hydrogen diffusion is completely suppressed when the photogenerated carriers are extracted from the diffusion region by the reverse electric field.

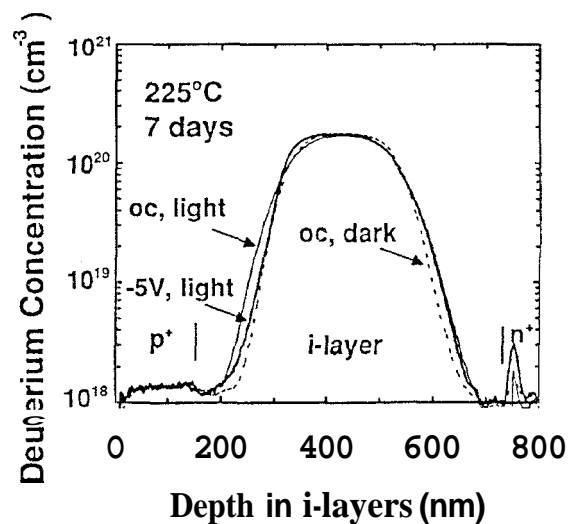


Figure 4: Deuterium profiles in p-i-n photodiodes annealed at 225°C for 7 days under open circuit (oc) conditions (dashed line), under oc and illumination (thin solid line), and under illumination and a reverse bias of 5V (thick line). The top chromium grid and the indium-tin-oxide layers were chemically etched before the SIMS analysis. p^+ and n^+ indicate the position of the corresponding doped layers.

IV. Discussion and Conclusions

The results presented in the previous section indicate that hydrogen diffusion and metastable defect formation in a-Si:H have many aspects in common. The most important is the fact that both processes are mediated by electronic carriers (i.e., electrons and/or holes)

and are accelerated when the carrier density is increased by illumination. From the comparison between diffusion experiments from a plasma and from a deuterated layer (see Fig. 2) we concluded that the main effect of the carriers is to increase the release rate of hydrogen atoms from Si-H bonds^[12]. If dangling-bonds are created following hydrogen release, a direct connection is established between hydrogen motion and metastable defect formation.

Two mechanisms can account for the interaction between carrier, hydrogen atoms, and defects. In the first the recombination of electron-hole pairs provide the energy necessary for hydrogen emission from a Si-H bonds and consequent defect formation^[16]. The second possibility is that the presence of carriers weaken Si-Si or a Si-H bonds, thereby decreasing the effective energy for hydrogen emission^[7].

Diffusion experiments in doped layers give additional support for a hydrogen-related metastability model. The relative changes in the defect density due to illumination are considerably smaller in doped and in compensated materials than in undoped a-Si:H.^[17,18] Similar observations apply to hydrogen diffusion in doped films. Finally, light-degradation experiments in solar cells indicate that the defect formation rate is suppressed under reversed biased solar cells^[19]. As is shown in Fig. 4, the hydrogen migration is also strongly suppressed under reverse bias.

These experimental results strongly support a close connection between hydrogen diffusion and light-induced defect generation and, therefore, an intrinsic, hydrogen-related mechanism for metastability in a-Si:H. Metastable defect formation under illumination is attributed to the increased rate of release of hydrogen from Si-H bonds in the presence of photogenerated carriers. According to this model, the efforts to improve material stability against defect formation should concentrate in obtaining a-Si:H films with more stable Si-H bonds.

References

1. D. L. Staebler and C. R. Wronski, *Appl. Phys. Lett.* **31**, 292 (1977).
2. S. Guha, J. Yang, W. Czubytyj, S. J. Hudgens, and M. Hack, *Appl. Phys. Lett.* **42**, 588 (1983); H. Pfeleiderer, W. Kusian, and W. Krueger, *Solid State Comm.* **49**, 493 (1984); W. den Boer, M. J. Geerts, M. Ondris, and H. M. Wenlinck, *J. Non-Cryst. Solids* **66**, 363 (1984).
3. T. B. Jackson and J. Kakalios, in *Amorphous Silicon and Related Materials*, Ed. by H. Fritzsche (World Scientific, Singapore, 1988), p. 247.
4. See, for instance, *Amorphous Silicon Materials and Solar Cells*, AIP Conf. Proc. **234**, Ed. by B. L. Stafford, (American Institute of Physics, New York, 1991).
5. D. Redfield and R. Bube, *Phys. Rev. Lett.* **65**, 464 (1990).
6. R. A. Street, *Phys. Rev. Lett.* **49**, 1187 (1982).
7. S. B. Zhang, W. B. Jackson, and D. J. Chadi, *Phys. Rev. Lett.* **65**, 2575 (1990).
8. M. Stutzmann, *Phil. Mag.* **B56**, 63 (1987).
9. H. Derschi, J. Stuke, and J. Beichler, *Appl. Phys. Lett.* **38**, 456 (1980).
10. D. E. Carlson and C. W. Magee, *Appl. Phys. Lett.* **33**, 81 (1978).
11. P. V. Santos and N. M. Johnson, *Appl. Phys. Lett.* **53**, 1181 (1993).
12. P. V. Santos, N. M. Johnson, and R. A. Street, *Phys. Rev. Lett.* **67**, 2686 (1991).
13. P. V. Santos and W. B. Jackson, *Phys. Rev.* **B46**, 4595 (1992).
14. H. M. Branz, S. E. Asher, and B. P. Nelson, in *Amorphous Silicon Technology-1992 Mat. Res. Symp.*, Vol. **219**, Ed. by A. Madam, Y. Hamakawa, M. J. Thompson, P. C. Taylor, and P. G. LeComber, (Material Research Society, Pittsburgh, 1992), p. 431.
15. P. V. Santos, N. M. Johnson, R. A. Street, M. Hack, R. Thompson, and C. C. Tsai, in print.
16. M. Stutzmann, W. B. Jackson, and C. C. Tsai, *Phys. Rev.* **B32**, 23 (1985).
17. A. Skumanich, N. M. Amer and W. B. Jackson, *Phys. Rev.* **B31**, 2263 (1985).
18. M. Stutzmann, *Appl. Phys. Lett.* **56**, 2313 (1990).
19. D. L. Staebler, R. S. Crandall, and R. Williams, *Appl. Phys. Lett.* **39**, 733 (1981).