

## Spin freezing of the spin-glass $\text{Fe}_{0.25}\text{Zn}_{0.75}\text{F}_2$ : a Mössbauer study

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**Abstract** We report  $^{57}\text{Fe}$  Mössbauer measurements in the diluted Ising antiferromagnet  $\text{Fe}_{0.25}\text{Zn}_{0.75}\text{F}_2$  at temperatures between 4.2 and 28 K. DC susceptibility measurements in the same sample show a spin-glass phase at a freezing temperature  $T_g = 10\text{K}$ . We found that for this concentration there is a competitive coexistence of spin-glass behaviour and antiferromagnetic order. The freezing temperature is frequency-dependent and follows a power law.

The phase transition behaviour of random magnetic systems continues to attract a great deal of attention both theoretically and experimentally. Very recently it has been shown that an antiferromagnet with intra-sublattice frustration shows a rich (H-T) phase diagram<sup>1</sup> (H=magnetic field, T=temperature). For H=0, a coexistence of antiferromagnetic (AF) and spin glass (SG) phases is observed. As H increases the SG phase is dominant. It has been shown recently<sup>2</sup> that the diluted Ising AF  $\text{Fe}_x\text{Zn}_{1-x}\text{F}_2$  with  $x \leq 0.31$  displays many features of a three-dimensional short-ranged SG system. At high Fe concentrations ( $x > 0.4$ ), under a uniform external magnetic field, this system is recognized as an excellent experimental realization of a  $d = 3$  random-field Ising model (RFIM) system<sup>3</sup>. DC susceptibility measurements for samples with concentrations  $x \leq 0.25$  reveal a SG phase appearing at low temperatures<sup>2</sup>. In particular for  $x = 0.25$ , the irreversibility temperature

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is at  $T_g = 10\text{K}$ . For this same sample ( $x = 0.25$ ) AC-susceptibility measurements confirm the features of SG behaviour: it has a peak in  $T$  whose amplitude decreases and which shifts to higher  $T$  as frequencies increase; as  $H_0$ , an external applied field parallel to the AC field, increases, the peak is reduced but does not shift in  $T$ <sup>4</sup>. By Mössbauer spectroscopy measurements in the same sample we have observed that a competitive coexistence of a SG phase and an AF order occurs below  $21\text{K}$ <sup>5</sup>.

In this paper, we report more detailed low-temperature Mossbauer measurements on this system with  $x = 0.25$ , and an analysis of the frequency dependence of the freezing temperature.

Figure 1 shows the Mössbauer spectra at 28, 21, 12 and 4.2 K. At 21 K we observe a superposition of a magnetic hyperfine spectra attributed to an SG phase and a doublet with broad linewidths which correspond to short-range AF order, as was shown previously<sup>5</sup>. Above 28 K we observe a paramagnetic quadrupolar spectrum with narrow linewidths. For  $T \leq 28\text{K}$  the spectra show an asymmetrical doublet with broad linewidth. At 4.2 K the spectrum is similar to that of pure  $\text{FeF}_2$ , but with a smaller hyperfine magnetic field ( $H_{hf}$ ) and larger linewidths due to the Zn dilution<sup>6</sup>. This is an indication that the  $H_{hf}$  stays perpendicular to the principal axis of the electric field gradient ( $EFQ$ ). For  $4.2\text{K} \leq T \leq 21\text{K}$  the spectra were fitted assuming a superposition of two static and gaussian distributions of hyperfine fields and diagonalizing the complete Hamiltonian. The  $P(H)$  distribution is shown in the right part of figure 1. This gaussian behaviour of  $P(H)$  is in accordance with Monte Carlo simulations<sup>7</sup> and theoretical results<sup>8</sup>. The introduction of a linear correlation between  $H_{hf}$  and the quadrupolar splitting ( $\Delta E_Q$ ) does not improve the fit and we obtain a very small correlation factor. The introduction of the  $EFQ$  asymmetry parameter  $\eta$  (value between 0.6 and 0.8) is necessary in order to obtain good fits.

According to the AC-susceptibility ( $\chi_{ac}$ ) measurements, the cusp temperature depends on the measuring frequencies  $f$ <sup>4</sup>. Many researchers have attempted to explain this phenomenon by introducing a time relaxation distribution for the spin systems. As the temperature decreases the relaxation time  $\tau$  begins to distribute

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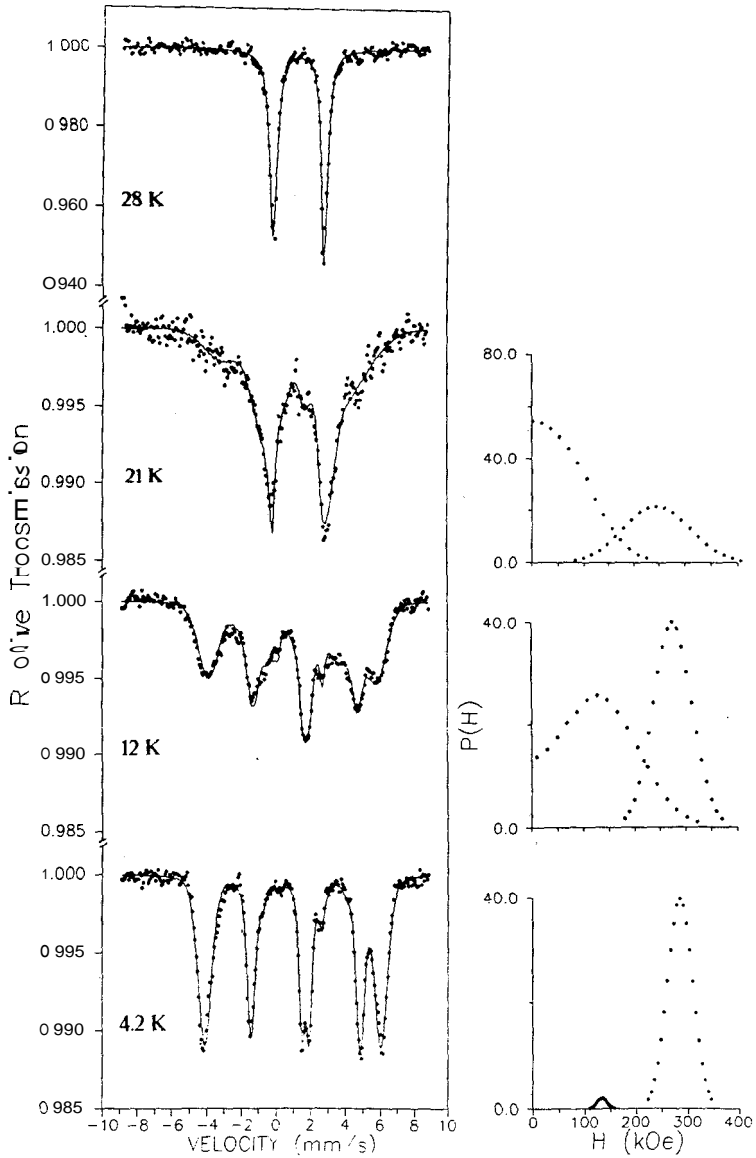


Fig. 1 -  $^{57}Fe$  Mössbauer spectra at 28, 21, 12 and 4,2K in  $Fe_{0.25}Zn_{0.75}F_2$ . Full lines are fittings with two static, gaussian distributions of  $H_{hf}$ . The P(H) curve is shown on the right of the corresponding spectrum.

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widely, and the distribution moves toward longer times and spreads broadly. Below the temperature  $Tg(f)$  at which a characteristic value  $r$ , of the relaxation time becomes equal to  $f^{-1}$ , only magnetic moments with  $r \leq \tau_c$  contribute to AC-susceptibilities, and  $\chi_{ac}$  decreases below  $Tg(f)$ . Thus we can recognize  $Tg(f)$  as a temperature where AC-susceptibility shows a maximum. Aruga et al.<sup>9</sup> have shown that the frequency law explaining the dependence of  $Tg(f)$  is of the form

$$f = f_0[(Tg(f) - Tg(0)/Tg(f)]^{z\nu} \quad (1)$$

where  $z$  is the dynamic exponent and  $\nu$  the **critical** exponent of the correlation length.

Mössbauer spectroscopy gives information concerning a single spin time-averaged within the nuclear Larmor precession time  $\tau_L$  of  $^{57}\text{Fe}$ . In the case of  $\text{Fe}_{0.25}\text{Zn}_{0.75}\text{F}_2$   $\tau_L$  is estimated to be  $1.6 \times 10^{-7}\text{s}$ , and  $f \cong 6.1 \times 10^6$  Hz. By Mössbauer spectroscopy the transition occurs at  $T = 21\text{K}$ . As the temperature decreases, spin fluctuations slow down gradually. Around  $21\text{K}$ , the fluctuation time of some spins becomes comparable to  $\tau_L$  and the spins appear to freeze. Thus we can consider the temperature of  $21\text{K}$  as  $Tg(f)$  for  $f \cong 6.1 \times 10^6$  Hz. Using the values of  $Tg(f)$  obtained by AC-susceptibility<sup>4</sup> and Mössbauer measurements, we can make a plotting of  $Tg(f)$  versus  $f$  (Fig. 2). We can see that the power law (1) is obeyed. The values of  $z\nu$  and  $f_0$  are around 7 and  $9 \times 10^9$  Hz, respectively. The value of  $z\nu$  is in accordance with that obtained by Rezende et.al.<sup>4</sup>. This value is still appropriate for an Ising SG. Nevertheless, the value of  $f_0$  obtained in this work is very different from that obtained in  $\text{Fe}_{0.5}\text{Mn}_{0.5}\text{TiO}_3$  by Aruga et. al.<sup>9</sup>, where  $f_0 = 1.4 \times 10^{13}$  Hz. This disagreement can be related to the AF phase present in  $\text{Fe}_{0.25}\text{Zn}_{0.75}\text{F}_2$ .

### **Acknowledgments**

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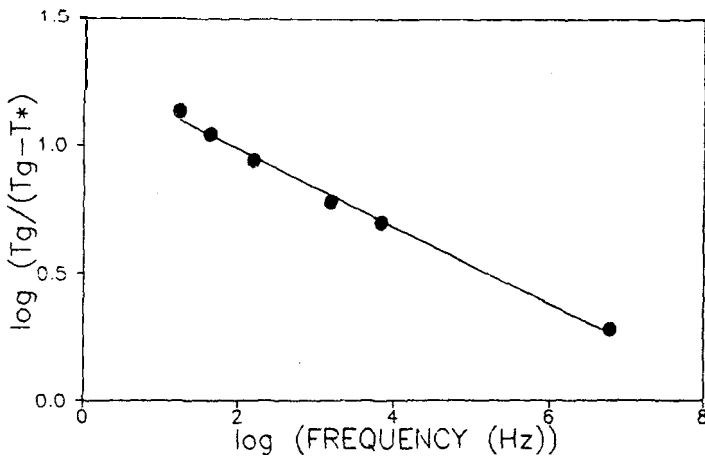


Fig. 2 - Log-Log plot of  $Tg(f)/(Tg(f) - Tg(0))$  versus frequency.

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### **Resumo**

Apresentamos medidas M6ssbauer de  $^{57}\text{Fe}$  no antiferromagneto diluido  $\text{Fe}_{0,25}\text{Zn}_{0,75}\text{F}_2$  em temperaturas entre 4,2 e 28 K. Medidas de Susceptibilidade magn6tica DC na mesma amostra mostram a exist6ncia de uma fase vidro de spin numa temperatura  $T_g = 10$  K. Para esta concentra76o mostramos que h6 uma coexist6ncia competitiva entre um comportamento vidro de spin e uma ordem antiferromagn6tica. A temperatura de congelamento 6 dependente da frequ6ncia e obedece a uma lei de pot6ncia.