

## Curvature Copies in Metric and Semi-Metric Manifolds

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Recebido em 18 de agosto de 1988

**Abstract** We present examples of curvature copies in metric and semi-metric spaces. These mappings involve geometrical objects of the manifolds and have some physical **implications** in the context of gravitational theories.

### ■ INTRODUCTION

In gravitational theories based on curved space-time the field equations are obtained from invariants of the curvature tensor **and/or** its contractions (Ricci tensor and scalar curvature). The usual arena for these theories is a metric **4-dimensional** manifold provided with a linear connection (Riemann-Cartan space  $U_4$ ) which is called Riemann space ( $V_4$ ) when the connection is symmetric. In this latter case the most remarkable theory of gravitation is, of course, General **Relativity**.

An interesting question concerning these geometrical structures is the **possibility** of the **existence** of different connections corresponding to the same curvature tensor. Thus different connections can generate the same curvature invariants which determine the field equations. Connections with this property **will** be called curvature copies here, in analogy with the field strength copies in non-abelian gauge theories<sup>1</sup>. **Field** copies have been studied also for the case of **internal** curvatures related to internal gauge transformations in the framework of Einstein theory of gravitation<sup>2</sup>.

The investigation of non-metric space-times has attracted less attention in recent years, since theories based on such manifolds **usually** reveal scarce **physical** meaning. In spite of this, some aspects of **invariance** properties and curvature copies in non-metric spaces **seem** to **deserve** consideration. The **projective**, or **Einstein  $\lambda$ -transformation**<sup>3</sup>, for instance, can be used in the construction of metrizable theories of gravitation<sup>4</sup>, or can be theories of gravitation<sup>4</sup>, or can be associated

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Work **partially** supported by **CNPq** (Brazilian Government Agency)

with gauge mappings for Yang-Mills fields<sup>6</sup>. The relation between **projective mappings** and conservation laws in quantum physics<sup>6</sup> has been **suggested as well**.

**In** this note we shall present some new examples of copies **in** the context of metric and semi-metric n-dimensional spocetime. The applications to 4-dimensional gravity may be specially interesting. We shall **follow** here the **definitions** and notations of **ref.7**.

## 2. PRELIMINARIES

Let us consider a general n-dimensional manifold **L**, **in** which exist: a rotation curvature, a torsion field and a curvature of segmentation. Let  $\Gamma_{\alpha\beta}^{\gamma}$  be the **affine** connection of **L<sub>r</sub>**. By adding a tensor  $P_{\alpha\beta}^{\gamma}$  to  $\Gamma_{\alpha\beta}^{\gamma}$  we obtain an object that transforms like a connection. This new connection

$$\bar{\Gamma}_{\alpha\beta}^{\gamma} = \Gamma_{\alpha\beta}^{\gamma} + P_{\alpha\beta}^{\gamma} \tag{1}$$

belongs to a manifold  $\bar{L}_n$ . From  $\bar{\Gamma}_{\alpha\beta}^{\gamma}$  we derive the curvature  $\bar{R}_{\alpha\beta\gamma}^{\delta}$  of  $\bar{L}_n$ . Thus we can impose the condition of invariance of curvature'

$$\bar{R}_{\alpha\beta\gamma}^{\delta}(\Gamma) = R_{\alpha\beta\gamma}^{\delta}(\Gamma) \tag{2}$$

for  $\bar{L}_n$  and  $L_n$ .

Def. 2.1 A set of connections in  $\bar{L}_n$  constructed by the addition of a tensor of third rank  $P_{\alpha\beta}^{\gamma}$  to a connection  $\Gamma_{\alpha\beta}^{\gamma}$  in  $L_r$  and satisfying the condition of invariance of curvature for  $\bar{L}_n$  and  $L_r$  (ref. eq.(2)) **will be called** a set of curvature field copies of  $\Gamma_{\alpha\beta}^{\gamma}$  in  $\bar{L}_n$ . The tensor  $P_{\alpha\beta}^{\gamma}$  **will be named** P-field.

From eq.(2) we get for the P-field the condition

$$\nabla_{[\alpha} P_{\beta]\gamma}^{\delta} + S_{\alpha\beta}^{\lambda} + P_{[\beta|\gamma]}.^{\lambda} P_{\alpha]\lambda}^{\delta} = 0 \tag{3}$$

where

$$S_{\alpha\beta}^{\lambda} \equiv \Gamma_{[\alpha\beta]}^{\lambda} \quad \text{and} \quad \nabla_{\nu}$$

is the covariant derivative associated with the connection  $\Gamma_{\alpha\beta}^{\lambda}$ .

We shall analyse specific examples of copies for which the P-field **is** generated by the mathematical structure of the initial (base) manifold. that **is**, copies which

involve quantities as the non-metricity tensor  $Q_{\alpha\beta\gamma} = -\nabla_{\alpha}g_{\beta\gamma}$  and the torsion tensor  $S_{\alpha\beta\gamma}$  and their respective contractions. We notice that this requirement is not necessarily fulfilled, for example, in the case of general projective copies

$$\bar{\Gamma}_{\alpha\beta}^{\lambda} = \Gamma_{\alpha\beta}^{\lambda} + \delta_{\beta}^{\lambda}\psi_{\alpha} \quad (4)$$

for an arbitrary vector field  $\psi_{\alpha}$ .

Def. 2.2 A curvature field copy which is generated by tensorial objects of the manifold [torsion and non-metricity tensors and their contractions] is called here an intrinsic geometrical curvature copy (IGCC, for short).

### 3. CURVATURE COPIES

As a first example we take a tensor in a semi-metric manifold  $L$ ,

$$P_{\alpha\beta}^{\lambda} = 2a g_{\alpha[\beta}Q_{\sigma]}g^{\sigma\lambda} + b Q_{\alpha}\delta_{\beta}^{\lambda} + c S_{\alpha}\delta_{\beta}^{\lambda} \quad (5)$$

where  $a, b, c$  are real numbers and

$$Q_{\gamma}g_{\alpha\beta} = Q_{\gamma\alpha\beta} \quad (\text{semi-metricity condition}), \quad (6)$$

$$S_{\alpha} = \frac{2}{n-1}S_{\alpha}^{\beta}. \quad (\text{torsion vector}). \quad (7)$$

Let us also assume the following restrictions for  $S$  and  $Q_{\alpha}$  in eq.(5)

$$\nabla_{\beta}Q^{\alpha} = 0 \quad (8)$$

$$Q_{\alpha}Q^{\alpha} = 0 \quad (9)$$

$$S_{\alpha\beta\gamma} = (a-1)g_{\gamma[\alpha}Q_{\beta]} \quad (10)$$

Theor. 3.1. - The P-field expressed by eq.(5) is a P-field of an IGCC if the conditions (8) - (10) are satisfied, with  $Q$  and  $S$  defined by eqs. (6) and (7), respectively.

The proof of this theorem **follows immediately** from the substitution of eq.(5) into eq.(3), with the conditions (8) - (10). We remark that in this case the original manifold  $L$  is semi-metric with non-symmetric connection, the torsion being determined by the null constant vector  $Q_\alpha$ , which expresses the existence of a curvature of segmentation.

Another **example** of a semi-metric curvature copy is provided by the  $P$ -field

$$P_{\alpha\beta}{}^\gamma = g_{\alpha\beta}Q^\gamma \quad (11)$$

with the assumptions

$$\nabla_\beta Q^\alpha = 0 = S_{\alpha\beta\gamma} \quad (12)$$

Theor. 3.2. - The  $P$ -field given by eq.(11) corresponds to an IGCC if the condition eqs.(6) and (12) are imposed.

Again the proof **follows** from direct substitution in to eq.(3). Here we have a copy in Weyl's geometry  $W_n$ , since

$$\nabla_\gamma g_{\alpha\beta} = -Q_\gamma g_{\alpha\beta}$$

and

$$S_{\alpha\beta\gamma} = 0$$

By considering now the  $P$ -field

$$P_{\alpha\beta}{}^\gamma = \delta_\alpha^\gamma S_\beta \quad (13)$$

with

$$S_{\alpha\beta\gamma} = S_{[\alpha}g_{\beta]\gamma} \quad (\text{vectorial irreducible component of the torsion}) \quad (14)$$

and

$$\nabla_\alpha S_\beta = 0 \quad (15)$$

we obtain a different copy in a semi-metric  $L$ ,

Theor. 3.3. - The P-field introduced by eq.(13) is connected with an IGCC if the assumptions eqs. (6), (14) and (15) are taken into account.

This follows by using the same substitutions referred to above.

We shall hereafter analyse the case of metric n-dimensional spaces, in which

$$\nabla_\gamma g_{\alpha\beta} = 0 \quad (\text{metricity condition}) \tag{16}$$

We point out that in this situation the imposition of a metricity condition for the copy manifold  $\bar{L}_n$ , as well as for the base manifold  $L_n$  introduces a symmetry restriction for the corresponding P-field.

Lemma 3.1. - In a curvature copy mapping of a metric manifold  $M_n$  into a metric manifold  $\bar{M}_n$  the P-field has the symmetry property  $P_{\alpha\beta\gamma} = P_{\alpha[\beta\gamma]}$ .

The proof follows from

$$\bar{\nabla}_\gamma g_{\alpha\beta} = \nabla_\gamma g_{\alpha\beta} = 0 \tag{17}$$

where  $\bar{\nabla}_\lambda$  is associated with the connection  $\bar{\Gamma}_{\alpha\beta}^\gamma$  in  $M_n$ .

Incidentally we also observe that if in this last case we further make the restriction  $S_{\alpha\beta\gamma} = 0$ , reducing the manifold to a Riemann space  $V_n$  then the P-field vanishes.

Theor. 3.4. - A curvature copy mapping of a Riemmanian n-dimensional manifold into itself is necessarily the identity mapping.

To prove this theorem we use the Lemma and the fact that the connection  $\Gamma_{\alpha\beta}^\gamma$  and its copy are symmetric. Thus in this case  $P_{\alpha\beta\gamma} = P_{(\alpha\beta)\gamma}$ , and once we have from the Lemma  $P_{\alpha\beta\gamma} = P_{\alpha[\beta\gamma]}$ , we are led to  $P_{\alpha\beta\gamma} = 0$ , which implies  $\bar{\Gamma}_{\alpha\beta}^\gamma = \Gamma_{\alpha\beta}^\gamma$ , i.e., the identity copy mapping.

We take now as an example of P-field in a metric n-dimensional space with torsion  $(U_n)$  the tensor

$$P_{\alpha\beta\gamma} = 2 g^{\lambda\gamma} g_{\alpha[\lambda} S_{\beta]} \tag{18}$$

with  $S_\lambda$  defined by eq.(7), plus the condition

$$S_\alpha S^\alpha = 0 \tag{19}$$

which gives, in view of the metricity,

$$\nabla_{\beta} S_{\alpha} = 0.$$

Theor. 3.5. - The P-field in eq.(18) is related to an IGCC if we impose eqs. (16) and (19).

This particular copy has an specialization, **derived** by imposing on the vielbein field  $e_{\alpha}^j$  in  $U_n$  the condition

$$\nabla_{\beta} e_{\alpha}^j = 0 \tag{20}$$

where  $j = 0, \dots, n - 1$  is an **internal** index. From eq.(20) we have

$$\Gamma_{\alpha}^j = e_j^{\gamma} \partial_{\alpha} e_{\beta}^j \tag{21}$$

which is the connection of the Weitzenbock space  $A_n$ . Relations (18) and (21) lead to

$$\Gamma_{\alpha\beta}^{\gamma} = e_j^{\gamma} \partial_{\alpha} e_{\beta}^j + \frac{2}{n-1} (\delta_{\beta}^{\sigma} \delta_{\alpha}^{\gamma} - g_{\alpha\beta} g^{\sigma\gamma}) e_j^{\lambda} \partial_{[\sigma} e_{\lambda]}^j \tag{22}$$

This copy is entirely determined by the vielbein **parallel** fields of  $A_n$ . We notice that not only the  $U_4$  space has been used in the formulation of alternative theories of the gravitational field, but also the 4-dimensional Weitzenbock space  $A_4$  has been adopted in this context<sup>8</sup>. In  $A_n$  the curvature tensor vanishes, so that, **according to eq.(2)**, the curvature of the copy space  $\bar{A}_n$  **also** vanishes.

Another subspace of  $U_n$ , besides  $A_n$  is the Riemann space  $V_n$ , which results from the assumption  $S_{\alpha\beta\gamma} = 0$  in  $U_n$ .

Copies in  $U_4$ , **such** as that given by eq. (18). with  $n = 4$ , **rise** the question of gravitational theories with identical free field equations but different couplings with matter. We can take, for instance, the interaction of spin  $-1/2$  particle with the **gravitational** field described by<sup>9</sup>

$$\mathcal{L}_M = \frac{i}{2} e_j^{\alpha} (\bar{\Psi} \gamma^j D_{\alpha} - \Psi \gamma^j D_{\alpha} \bar{\Psi}) - m \Psi \bar{\Psi} \tag{23}$$

where

$$D_{\alpha} = \partial_{\alpha} - \frac{i}{2} A_{ij\alpha} S^{ij}$$

$$S^{ij} = \frac{i}{4} [\gamma^i, \gamma^j]$$

is the infinitesimal generator of the **internal** Lorentz group:  $A_{ij\mu}$  is the Lorentz gauge field associated with the **affine** connection of  $U_4$ :  $\gamma^i, \gamma^j = 2\eta^{ij}; j = 0, \dots, 3$ . The copy  $\bar{\Gamma}_{\alpha\beta}^\gamma$  originates a field  $\bar{A}_{ij\mu} \neq A_{ij\mu}$ , and consequently a modified matter Lagrangian  $\bar{\mathcal{L}}_M$ . In the **4-dimensional teleparallel** space  $A_4$  the internal connection  $A_{ij\mu}$  vanishes, so that in the minimal coupling **instance** we have  $\bar{\mathcal{L}}_M = \mathcal{L}_M$ , and the copy does not alter the interaction. However, **if** we take a non-minimal coupling Lagrangian density like

$$\mathcal{L}'_M = iK\sqrt{-g}\Psi\gamma^k\Psi e_k^\mu S_m u \tag{24}$$

where  $K$  is a constant and

$$S_\mu = \frac{2}{3} e_k^\mu \partial[\mu e^k]$$

we obtain  $\bar{\mathcal{L}}'_M = 2\mathcal{L}'_M$ , since in the copy manifold  $\bar{S}_\mu = 2S_\mu$ . Thus even in the  $A_4$  case the gravitational interaction of fermions may be affected by the copies.

Finally we want to mention the effect of those mappings on trajectories in  $U_4$ . Autoparallel **curves**<sup>10</sup> in the copy space differ in general from those in the original one, even if the gravitational equations do coincide. This deviation depends, of course, on the symmetric part of the field copy, since the symmetric part of the connection appears in the equation that defines autoparallel. Extremal curves are derived directly from the metric and do not undergo any copy effect. Concerning the motion of spin test particles outside of the matter, it is reasonable to expect some **influence** of copies on the trajectories in  $U_4$  theories which have propagation of torsion.

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### Resumo

Apresentamos exemplos de cópias de curvatura em espaços métricos e semi-métricos. Tais mapeamentos envolvem objetos geométricos das variedades, e têm implicações físicas no contexto de teorias gravitacionais.