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New Algebraic Tables of SU(2) Quantitieo

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Abstract Formulas for Clebsch-Gordan Coefficients, 6-j and 9-j symbols of SU(2) are presented in a ready-to-program way for obtaining algebraic tables. An excerpt of the complete tables is also presented.

1. INTRODUCTION

Formulas for SU(2) Clebsch-Gordan Coefficients, 6-j and 9-jsymbols can be found in any book on application of Group Theory to Quantum Mechanics¹⁻³. These formulas can be easily programmed to produce *numerical* values. From them can also be obtained formulas with the simple structure of a polynomial multiplied by the square root of the ratio of two polynomials. We do that in the following sections.

Algebraic formulas of the SU(2) quantities with the structure mentioned above, besides having an interest by their own, are useful in several instances. With one single such algebraic formula one obtains not only one but an infinity of numerical values right away by using a pocket calculator (or even by hand calculation) just by ascribing numerical values to the variables. The formulas for the analogous quantities for higher dimensional unitary and orthogonal groups⁴ have sometimes the SU(2) quantities as internal blocks. Hence, simpler algebraic formulas for the previous would require simpler algebraic formulas for the latter. These simpler algebraic formulas for higher dimensional unitary and orthogonal groups are also relevant in the development of the so called Mlcroscopic Model of the Collective Nuclear Motion⁵. This

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was in fact our main motivation for the present work. Even the numerical evaluation of the SU(2) quantities will be benifited by these simpler algebraic formulas in the numerical precision as well in the computing time.

Another possible application of these formulas 1s the study of the non-trivial zeros of the Racah (6-j) cuefficients⁶.

The **most** extensive algebraic tables of these **SU(2)** quantities have been published by D.A. Varchalovitch, A. N. Moskalev and V.K. Khersonski**i⁷.** Our tables are an independent check^e and an extension of their tables.

2. CLEBSCH-GORDAN COEFFICIENTS

For the SU(2) Clebsch-Gordan **Coefficients** we start from well known Racah's formula¹⁻³ $C_{m_{1}m_{2}m}^{j_{1}j_{2}j} = \delta(m_{1}+m_{2},m) \frac{[(2j+1)(j_{1}+j_{2}-j)!(j_{1}-j_{2}+j)!(-j_{1}+j_{2}+j)!(j+m)!(j-m)!]}{(j_{1}+j_{2}+j+1)!(j_{1}+m_{1})!(j_{1}-m_{1})!(j_{2}+m_{2})!(j_{2}-m_{2})!} \int_{-1}^{1} \frac{(j_{1}+j_{2}+j)!(j+m)!(j-m)!}{(j_{1}+j_{2}+j+1)!(j_{1}+m_{1})!(j_{1}-m_{1})!(j_{2}+m_{2})!(j_{2}-m_{2})!} \int_{-1}^{1} \frac{(j_{1}+j_{2}+j)!(j+m)!(j-m)!(j-m)!}{(j_{1}+j_{2}+j+1)!(j_{1}-m_{1})!(j_{1}-m_{1})!(j_{2}+m_{2})!(j_{2}-m_{2})!} \int_{-1}^{1} \frac{(j_{1}+j_{2}+j)!(j+m)!(j-m)$

$$\times \sum_{z} \frac{2}{z! (j+m-z)! (j-j_1+j_2-z)! (j_1-j_2-m+z)!}$$
(2.1)

In order to obtain formulas with the structure mentioned in the introduction, like those given in Condon-Shortley's tables, one must sacrifice the generality of eq. (2.1) by assigning numerical values to some of its entries. The procedure is better seen if we re-write eq. (2.1) putting $j = j_1 + a$, $j_2 = b$; $m_2 = a$, thus obtaining

$$C_{m-\alpha \ \alpha \ m}^{j_{1} \ b \ j_{1}+a} = \frac{\left[\frac{(2j_{1}+2a+1)(b-a)!(2j_{1}+a-b)!(b+a)!(j_{1}+m+a)!(j_{1}-m+a)!}{(2j_{1}+b+a+1)!(j_{1}+m-\alpha)!(j_{1}-m+\alpha)!(b+\alpha)!(b-\alpha)!}\right]^{\frac{1}{2}}{\left[\frac{(-)^{b+\alpha+z}(j_{1}+m+a+b-\alpha-z)!(j_{1}-m+\alpha+z)!}{z!(j_{1}+m+a-z)!(a+b-z)!(j_{1}-m-b+z)!}\right]}$$
(2.2)

To be able to perform the summation, the terms have to be expressed as ratios of polynomials. This is done by matching, for each factorial in numerator, a factorial in denominator such that the difference of their arguments be a numerical integer. From eq. (2.1) one sees that the only way to obtain this is by assigning numerical values to a, b and a.

By use of the following generalization of Pochhammer's symbol

$$P(f;k;\Delta) = f(f+\Delta)(f+2\Delta)...(f+(k-1)\Delta), \text{ for } k = \text{integer} > 0,$$

=1, for $k = 0,$ (2.3)

and the relations

 $b\pm \alpha = \text{integer} \ge 0$; $b\pm a \approx \text{integer} \ge 0$; $a+\alpha = \text{integer}$, (2.4) ensured by the branching rules, eq. (2.2) can be written as

$$c_{m-\alpha \ \alpha \ m}^{j_{1} \ b \ j_{1}+\alpha} - \left[\frac{(b-a)!(2j_{1}+2a+1)}{(b+\alpha)!(b-\alpha)!(b+a)!P(2j_{1}+a-b+1;2b+1;1)} \right]^{\frac{1}{2}} \\ \times \left\{ \begin{bmatrix} 2^{-(\mu-1)(a+\alpha)}P(\mu(j_{1}+m-\alpha+1);a+\alpha;\mu) \end{bmatrix}^{\frac{1}{2}} \\ \text{or} \\ \left[2^{(\mu-1)(a+\alpha)}P(\mu(j_{1}+m+a+1);-a-\alpha;\mu) \end{bmatrix}^{-\frac{1}{2}} \right\} \\ \times \left\{ \begin{bmatrix} 2^{-(\mu-1)(a-\alpha)}P(\mu(j_{1}-m+\alpha+1);a-\alpha;\mu) \end{bmatrix}^{\frac{1}{2}} \\ \text{or} \\ \left[2^{(\mu-1)(a-\alpha)}P(\mu(j_{1}-m+a+1);-a+\alpha;\mu) \end{bmatrix}^{-\frac{1}{2}} \right\} \\ \times \frac{1}{2^{2}(\mu-1)b} \sum_{z=0}^{a+b} (-)^{b+\alpha+z} \begin{bmatrix} a+b \\ z \end{bmatrix} P(\mu(j_{1}+m+a+1-z);b-\alpha;\mu) \\ \times P(\mu(j_{1}-m-b+1+z);b+\alpha;\mu) \\ \text{where} \\ \mu = 1 , \text{ for even } b$$

 $\mu = 1$, for even b $\mu = 2$, for odd b and in the round brackets one must take the terms that have P's with second argument positive.

All the numerical constants that appear in eq. (2.5) are integer, so it can be evaluated exactly using integer arithmetic. We have made a computer program to evaluate eq. (2.5) as an exact algebraic expression having j_1 and m as free variables and b, a and a as numerical values. This program produced a table of algebraic formulas for b = 1/2, 1, $3/2, \ldots, 8$ and $-b \le \alpha, \alpha \le b$. In table 1 we present part of this table. [The complete table, as well those for 6-j and 9-j symbols, are available upon request from the author J.A.C.A].

3. (6-j) SYMBOLS

Taking the formula¹⁻³

$$\begin{cases} j_{1}j_{2}j_{3} \\ \ell_{1}\ell_{2}\ell_{3} \\ \ell_{1}\ell_{2}\ell_{3} \end{cases} = (-1)^{j_{1}+j_{2}+\ell_{1}+\ell_{2}} \left[\frac{\left(-j_{1}+j_{2}+j_{3}\right) : \left(j_{1}+j_{2}-j_{3}\right) : \left(j_{1}-j_{2}+j_{3}\right) : \left(-j_{1}+\ell_{2}+\ell_{3}\right) : \left(j_{1}-\ell_{2}+\ell_{3}\right) : \left(j_{1}-\ell_{2}+\ell_{3}\right) : \left(j_{1}-\ell_{2}+\ell_{3}\right) : \left(j_{1}-\ell_{2}+\ell_{3}\right) : \left(j_{1}-\ell_{2}+\ell_{3}\right) : \left(\ell_{1}-\ell_{2}+\ell_{3}\right) : \left(\ell_{1}-\ell_{2}+\ell_{3}\right) : \left(\ell_{1}-\ell_{2}+\ell_{3}\right) : \left(\ell_{1}-\ell_{2}+\ell_{3}\right) : \left(\ell_{1}-\ell_{2}+\ell_{3}\right) : \left(\ell_{1}-\ell_{2}+\ell_{3}\right) : \left(\ell_{1}+\ell_{2}+\ell_{3}\right) : \left(\ell_{1}+\ell_{2}+\ell_{3}\right) : \left(\ell_{1}+\ell_{2}+\ell_{3}\right) : \left(\ell_{1}+\ell_{2}+\ell_{3}\right) : \left(\ell_{1}+\ell_{2}+\ell_{3}\right) : \left(\ell_{1}+\ell_{2}+\ell_{3}\right) : \left(\ell_{1}+\ell_{2}+\ell_{3}+1\right) : \left(\ell_{1}+\ell_{3}+\ell_{3}+\ell_{3}+1\right) : \left(\ell_{1}+\ell_{3}+\ell_{3}+\ell_{3}+\ell_{3}+1\right) : \left(\ell_{1}+\ell_{3}+\ell_{3}+\ell_{3}+\ell_{3}+1\right) : \left(\ell_{1}+\ell_{3}+\ell_{3}+\ell_{3}+\ell_{3}+\ell_{3}+\ell_{3}+1\right) : \left(\ell_{1}+\ell_{3}+\ell_{3}+\ell_{3}+\ell_{3}+\ell_{3}+1\right) : \left(\ell_{1}+\ell_$$

and putting

$$j_1 = a; \quad j_2 = b; \quad j_3 = c; \quad k_2 = c + a; \quad k_3 = b + \beta$$
 (3.2)

one obtains

$$\frac{a \ b \ c}{\ell_{R_{1}} \ c+a \ b+\beta} = (-)^{a+b+c+\ell_{1}+\alpha} \left[\frac{(-a+b+c)! \ (a-b+c)! \ (a+b-c)! \ (-a+b+c+\alpha+\beta)!}{(a+b+c+1)! \ (a+b+c+\alpha+\beta+1)!} \right]$$

$$\times \frac{(a+b-c-\alpha+\beta)! \ (a-b+c+\alpha-\beta)! \ (2b-\ell_{1}+\beta)! \ (\ell_{1}-\beta)! \ (2c-\ell_{1}+\alpha)! \ (\ell_{1}-\alpha)! \ (\ell_{1}-\alpha)! \ (\ell_{1}+\alpha)!}{(2b+\ell_{1}+\beta+1)! \ (2c+\ell_{1}+\alpha+1)!} \right]^{\frac{1}{2}}$$

$$\times \frac{(-)^{Z} \ (a+b+c+\ell_{1}+\alpha+1-z)!}{z! \ (a+b-c-z)! \ (\ell_{1}+\alpha-z)! \ (a-b+c+\alpha-\beta-z)! \ (\ell_{1}-\beta+z)! \ (-a+b+c-\ell_{1}+\beta+z)! \ (-\alpha+\beta+z)!}{(a+b+c+\ell_{1}+\alpha+1-z)!}$$

(3.3)

(3.1)

Table I - Algebraic expressions of Clebsch-Gordan Coefficients obtained from Eq.(2.5). For printing canvenience us put J in place of j and M in place of m.

$$\begin{bmatrix} J & 5/2 \ J^{+5/2} \\ H^{-1/2} \ 1/2 \ H \\ H^{-1/2} \ 1/2 \ H \\ = \begin{bmatrix} \overline{5(2J+2H+1)(2J+2H+3)(2J+2H+5)(2J-2H+5)(2J-2H+5)} \\ 1/2 \ 1/$$

As in the case of the Clebsch-Gordan coefficients, the generality must be sacrificed in order to have an expression with the desired simple structure. In this case the greatest generality is limited to 3 variables, We take a, b, c as free variables and ascribe numerical values to R, a and β . With this choice, there are in eq. (3.3) four numerical constraints on the values of z:

$$z \ge 0 ; \quad \ell_1 + \alpha - z \ge 0 ; \quad \ell_1 - \beta - z \ge 0; \quad -\alpha + \beta + z \ge 0$$
(3.4)

from which it follows that

$$z_{min} = \max(0, \alpha - \beta) ; \quad z_{max} = \min(\ell_1 + \alpha, \ell_1 - \beta) \quad (3.5)$$

There are then 4 possible ranges for z, according to the possibilities:

1)
$$\alpha + \beta \leq 0$$
 and $\alpha - \beta \leq 0$; 2) $a + \beta \leq 0$ and $a - 6 \geq 0$
3) $\alpha + \beta \geq 0$ and $\alpha - \beta \leq 0$; 4) $\alpha + \beta \geq 0$ and $\alpha - \beta \geq 0$
(3.6)

Taking into account eq. (3.6) one can re-express eq. (3.3) in terms of generalized **Pochhammer's** symbols as

$$\begin{cases} a & b & c \\ \ell_{1} & c+\alpha & b+\beta \end{cases} = \frac{(-)^{s}}{(\ell_{1}+\alpha)!} \frac{\left[\frac{P(n_{1};k_{1};1)P(s+n_{s};k_{s};1)P(s-2\alpha+n_{a};k_{a};1)}{P(n_{2};k_{2};1)P(2b-\ell_{1}+\beta+1;2\ell_{1}+1;1)} \right] \\ \times \frac{P(s-2b+n_{b};k_{b};1)P(s-2c+n_{a};k_{a};1)}{(2c-\ell_{1}+\alpha+1;2\ell_{1}+1;1)} \frac{1}{2} \times \sum_{\substack{z=z \\ z=z \\ m \mid n}}^{z} (-)^{\ell_{1}+\alpha+z} \binom{\ell_{1}+\alpha}{z} P(s+\bar{n}_{s};\bar{k}_{s}-z;1)$$

$$\times P(s-2a+\bar{n}_a+z;\bar{k}_a-z;1)P(s-2b+\bar{n}_b-z;\bar{k}_b+z;1)P(s-2c+\bar{n}_c-z;\bar{k}_c+z;1)$$

$$\times P(\ell_1 - \beta - z + 1; z; 1) P(-\alpha + \beta + z + 1; \ell_1 + \alpha - z; 1)$$
(3.7)

where

$$s = a + b + c ; \quad \bar{n}_{a} = -k_{1} + \beta + 1 ; \quad \bar{n}_{b} = a - \beta + 1$$

$$\bar{n}_{c} = 1 ; \quad k_{a} = k_{s} ; \quad k_{c} = k_{1} = k_{2} = k_{b} ; \quad \bar{k}_{c} = \bar{k}_{b}$$
(3.8)

and the other parameters are given in the table 2, according to the four possibilities of eq. (3.6).

Table 2 - Parameters of eq. (3.7) for the four possibilities given in eq. (3.6).

Case	ns	k _s	na	n _b	k _b	n	n ₁	n ₂	n _s	κ _s	\bar{k}_b
1	α+β+2	-α-β	α+β+1	α-β+1	-α+β	1	l_1-β+1	$\ell_1 + \alpha + 1$	2	l1+α	0
2	α+β+2	-α-β	α+β+1	1	α-β	-α+β+1	$l_1+\beta+1$	l,-α+1	2	ℓ 1+α	-α+β
3	2	α + β	1	α-β+1	-α+β	1	l_1-β+1	l l 1+α+1	α+β+2	ℓ 1-β	0
4	2	α+β	1	1	α-β	-α+β+1	ℓ1+β+1	$\ell_1^{-\alpha+1}$	α+β+2	ℓ 1 ^{-β}	-α+β

For tabulation purposes it is enough to consider one of the 4 cases of eq. (3.6), since the other cases can be obtained by using the symmetries of the 6-j symbols. For general use however it is better to make a single program that includes all the cases.

We have made also a computer program to evaluate eq.(3.7) as an exact algebraic expression having a.. b, c as free variables and R a and β as numerical values. This program produced a tableofalgebraic expressions for $\ell_1 = 1/2$, 1, $3/2, \ldots, 8$ and all allowed values of a and β such that $a + \beta \leq 0$ and $a - \beta \leq 0$ (case 1 of eq. (3.6)). In table 3 we reproduce part of this table.

4. (9-j) SYMBOLS

We start again from a well known formula $^{1-3}$

$$\begin{cases} a & b & c \\ \ell_1 & \ell_2 & \ell_3 \\ k_1 & k_2 & k_3 \end{cases} = \sum_{z} (-)^{2z} (2z+1) \left\{ \begin{matrix} a & k & z & k_2 & \ell_1 & z & \ell_3 & b & z \\ \ell_3 & b & c \end{matrix} \right\} \left\{ \begin{matrix} a & k_3 & k_1 & \ell_3 & k_2 & \ell_1 & \ell_2 \\ \ell_3 & b & c \end{matrix} \right\}$$
(4.1)

Taking l_1 , l_2 , l_3 as numerical constants, the branching rules require that k_1 , k, k_3 have the form

Table III-Algebraic expressions of $(\delta - j)$ symbols obtained from Eq.(3.8) The quantity X is defined as -a(a+1)+b(b+1)+c(c+1). For printing convenience we put capital letters A,B,C in place of low rase letters a,b,c and P(f;k) in place of P(f;k;1).

$$\begin{cases} A = C \\ S/2 = C-5/2 = B+1/2 \\ = (-)^{S} \left[\frac{10(S)(S+1)(S-2A-1)(S-2A)(S-2B-2)(S-2B-1)(S-2B)}{P(2E-1;6)} \right]^{1/2} \\ + \frac{(S-2B+1)(S-2C+2)(S-2C+3)}{P(2E-4;6)} \right]^{1/2} \\ \begin{cases} A = C \\ S = C-1 = S \\ = (-)^{S} \left[\frac{12(S+1)(S-2A)(S-2B)(S-2C+1)}{P(2E-2;7)P(2C-3;7)} \right]^{1/2} \\ + \frac{(S-2C+1)(S-2C+2)(S-2C+3)}{P(2E-4;6)} \right]^{1/2} \\ + \frac{(S-2C+3)}{S} \left[\frac{12(S+1)(S-2A)(S-2B)(S-2C+1)}{P(2E-2;7)P(2C-3;7)} \right]^{1/2} \\ + \frac{(S-2C+3)}{S} \left[\frac{12(S+1)(S-2A)(S-2B-2)(S-2B-1)(S-2B-1)(S-2C+1)(S-2C+2)}{P(2E-1;7)} \right]^{1/2} \\ + \frac{(S-2C+3)}{P(2C-4;7)} \right]^{1/2} (2(BC+B-2C-2)+3X) \\ \begin{cases} A = C \\ S = C-1 \\ S = C \\ S = C$$

$$k_1 = a + \alpha; \ k_2 = b + \beta; \ k_3 = c + \gamma$$
 (4.2)

with

$$-\ell_1 \leq \alpha \leq \ell_1 \; ; \; - \; \ell_2 \leq \beta \leq \ell_2 \; ; \; -\ell_3 \leq \gamma \leq \ell_3 \tag{4.3}$$

Therefore eq. (4.1) can be re-written as

$$\begin{cases} a & b & c \\ k_1 & k_2 & k_3 \\ a+\alpha & b+\beta & c+\gamma \end{cases} \begin{pmatrix} (-)^{2b} \sum_{\delta} (-)^{2\delta} (2b+2\delta+1) & \{a & b & c \\ k_3 & c+\gamma & b+\delta \end{cases} \\ \times \{ (c+\gamma) & a & (b+\delta) & (b+\delta) & (k_1+\lambda) & b \\ k_1 & (b+\delta)+(\beta-\delta) & a+\alpha \end{pmatrix} \{ (b+\delta) & (k_1+\lambda) & b \\ k_2 & b+\beta & (k_1+\lambda) - \lambda \end{cases}$$
(4.4)

where $X = \ell_3 - \ell_1$ and δ runs through $-\ell_3 \leq \delta \leq \ell_3$ subject to $-\ell_1 \leq (\beta - \delta) \leq \ell_1$. The entanglement of the (6-*j*) symbols in eq. (4.4) and the need of having ℓ_1 , α and β in eq. (3.7) as numerical values prevent us of obtaining general tables of (9-*j*) symbols in more than 3 free variables.

We have made a computer program to evaluate eq. (4.4) as an exact algebraic expression using a subroutine to evaluate the (6-j) symbols. This program produced a table of algebraic expressions for the (9-j) symbols, eq.(4.4), with $\ell_1 \leq \ell_2 \leq \ell_3 \leq 3$, for $\ell_1 = 1/2, 1, 3/2, 2, 5/2, 3$ and all porssibilities for a, β and γ . Using the symmetries of the (9-j) symbols, all the case with $\ell_2 \leq 3$ are covered by this table. In table 4 we present some items of the complete table.

In despite of the apparent simplicity of eq. (4.4), the coding of a program to evaluate it as an exact algebraic expression reaches practically the greatest degree of complexity in group theoretical algebraic calculations.

One of us (V.Vanagas) wishes to thank the Brazilian agency FAPESP for the finantial support of a six-month stay in the Instituto de Física Teorica, when this work was initiated. Table IV - Algebraic expressions of (9-J) symbols obtained from Eq.(4.4) The quantity Z is defined as a(a+1)+b(b+1)-c(c+1). For printing convenience we put capital letters A.B.C in place of low case letters a.b.C and P(f;k) in place of P(f;k;1).



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 We detected two errors in the 9-j tables of Ref.7. The correct ex-

pressions are given by the first **two** entries of our table 4.

Resumo

Apresentam-se fórmulas para Coeficientes de Clebsch - Gordan, símbolos 6-j e símbolos 9-j apropriadas para construção de tabelas algébricas via computador. Un resumo das tabelas completas obtidas dessas fórmulas é também apresentado.