

The Green's Function of a Relativistic Free-Electron in a Uniform Magnetic Field

J. BELLANDI FILHO and E. S. CAETANO NETO†

Instituto de Física Gleb Wataghin, UNICAMP, Campinas

and

S. M. L. PAVÃO

Instituto de Física Teórica, São Paulo

Recebido em 28 de Agosto de 1978

A closed representation is obtained for the Green's function of the Dirac equation of a free-electron in a uniform magnetic field.

Deduz-se uma representação fechada para a função de Green para a equação de Dirac para um elétron livre na presença de um campo magnético uniforme.

1. DERIVATION OF THE GREEN'S FUNCTION

The present note contains a simple derivation of the Green's function for a Dirac particle in a uniform magnetic field.

The Green's function satisfies the following inhomogeneous Dirac equation

$$(i\gamma_{\mu}\pi^{\mu} - \mu)S(x,y) = \delta^4(x-y) \quad (1)$$

where $\pi^{\mu} = \partial^{\mu} - eA^{\mu}$ is the energy-momentum four vector;

$$A = (0, -1/2 Hx, 1/2 Hx_1, 0)$$

† With a FAPESP - São Paulo - Brasil - Fellowship.

is the electromagnetic four vector potential; \vec{H} is the uniform magnetic field in the z -direction and μ is the electron mass.

We can make the following Ansatz

$$S(x, y) = (i\gamma_\mu \pi^\mu + \mu) F(x, y) \quad (2)$$

and as A does not depend on x_0 and x_3 , we may write $F(x, y)$ as a Fourier transform of a transversal Green's function

$$F(x, y) = \frac{1}{(2\pi)^2} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} dp_0 dp_3 \exp[i p_0 (x_0 - y_0) - i p_3 (x_3 - y_3)] K(\vec{s}, \vec{s}'; p_0, p_3) \quad (3)$$

which satisfies the following differential equation

$$[\vec{p}_0^2 - p_3^2 - 2\omega\sigma_3 - \mu^2 - \omega^2 s^2 + \nabla_\perp^2 - 2\omega L_3] K(\vec{s}, \vec{s}'; p_0, p_3) = \delta(\vec{s}' - \vec{s}) \quad (4)$$

with $\vec{s} \equiv (x_1, x_2)$ and $\omega = eH/2$. Assuming $K(\vec{s}, \vec{s}'; p_0, p_3)$ diagonal on σ_3 we may write $\sigma_3 K = sK$ with $s = \pm 1$.

As in the non-relativistic case (Bellandi, Zimmerman-1975) we can firstly seek for $K(\vec{s}, \vec{s}'; p_0, p_3)$ in the series form

$$K(\vec{s}, \vec{s}'; p_0, p_3) = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} K^m(\vec{s}, \vec{s}'; p_0, p_3) \exp[i m(\phi - \phi')] \quad (5)$$

where m are the eigenvalues of L_3 ; s and ϕ are the polar coordinates in the (x_1, x_2) plane.

The partial Green's function $K^m(s, s'; p, p)$ satisfies the following differential equation

$$\left[\frac{d^2}{ds^2} + \frac{1}{s} \frac{d}{ds} - \frac{m^2}{s^2} - \omega^2 s^2 + \lambda - 2m\omega \right] K^m(s, s'; p_0, p_3) = (ss')^{-1/2} \delta(s-s') \quad (6)$$

with $\lambda = p_0^2 - p_3^2 - 2\omega s - \mu^2$. A solution of this equation is (we put $\omega=1$)

$$K^m(s, s'; p_0, p_3) = \Gamma\left(\frac{1+|m|+m-\lambda/2}{2}\right) (ss')^{-1} M\left(\frac{\lambda-m}{2}, \frac{|m|}{2} (s<), \frac{\lambda-m}{2}, \frac{|m|}{2} (s>)\right). \quad (7)$$

Here $s <$ and $s >$ are the lesser and greater of s' and s' respectively and M and W are the Whittaker functions. Replacing the product of the Whittaker function in Eq.7 by an integral representation in the complex plane (Buchholz, 1969) we can perform the summation in Eq. 5 and the final expression for the transversal Green's function is

$$K(\vec{s}, \vec{s}'; p_0, p_3) = \frac{1}{2\pi} \Gamma\left(\frac{1}{2} - \frac{\lambda}{4}\right) \exp\left[-i \frac{|\vec{s} \times \vec{s}'|}{|\vec{s} - \vec{s}'|}\right] \frac{1}{|\vec{s} - \vec{s}'|} W_{\frac{\lambda}{4}}^{\lambda}; 0(|\vec{s} - \vec{s}'|^2).$$

(8)

Since this function is an analytical function, Eq. 8 is valid for arbitrary values of A . This expression must be interpreted as a 4×4 matrix, and for a defined value of s we must write $1/2(1 + s \Sigma^3)K$; $\Sigma^3 = \sigma^{12}$.

REFERENCES

- Bellandi Filho, J. and Zimmerman AH. 1975 *Lett.N.Cimento* 14, 521.
 Buchholz H. 1969 - The confluent Hypergeometric function, in *Springer Tracts in Natural Philosophy* Vol.15 (Berlin).