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Atomic Center Interactions in BaO; Al₂O₃; B₂O₃ Glasses Containing Silver*

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The EPR study of borate glasses, with 30% of Ba0 and 5% mole of silver, X-irradiated at 77° K. showed Ag⁰ and Ag⁺⁺ centers. In addition were detected the boron electron center (BEC) and the boron hole center. The silver centers and BEC were studied in detail and the spin Hamiltonian parameters are given. The different AR₂⁰ concentrations exerted only little influence on the tabulated constants. The Ag₂ center was not observed, indicating that the collisions between Ag atoms are notvery frequent in these glasses, even during the process of thermal bleaching. The hfs of the boron electron center sufferes severe variation as one compared the spectra of base glasses and those containing silver. This is an evidence of the action of silver as a glass modifier like the alkali atoms.

0 estudo de EPR dos vidros com boratos, contendo 30% de Ba0 e 5% mol de prata, irradiados com raios X a 77° K, revelou a presença dos centros Ag⁰ e Ag⁺⁺. Foram também observados o centro do elétron do boro (BEC) e o centro de buraco do boro. Os centros de prata e o EEC foram estudados em detalhe, sendo obtidos os parâmetros da Hamiltoniana de spin. As diferentes concentrações de AR 0 exerceram apenas uma pequena influência sobre as constantes expressas nas Tabelas. O centro de Ag⁺ não foi observado, indicando que as colisões entre os átomos de Ag não são muito frequentes nestes vidros, mesmo durante o processo de decaimento térmico. A constante de interação hiperfina do centro de elétron do boro sofreu acentuada variação quando da comparação dos espectros dos vidros-base

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com os que contêm prata. Este fato torna evidente a atuação da prata com modificadora do vidro, à semelhança dos átomos alcalinos.

1. INTRODUCTION

The study of **atomic** centers **in** glassy matrices has received increasing attention in the last decade.

In this work, we intend to add more data about the silver interactions in Ba0; $A_2^0 0_3$; $B_2^0 0_3$ glasses. The silver centers were already investigated in several amorphous matrices, being related to the EPR spectra of Ag^0 , Ag^{++} and Ag_2^+ centers. These centers were induced by X -irradiation at liquid nitrogen temperature. During irradiation, electrons and holes are liberated in the sample, both being trapped by the Ag^+ ion, which is present in the sample in its original state. The production of the centers is as follows:

$$Aa^{+} + e^{-} \rightarrow Ag^{0} , \qquad (1.1)$$

$$Ag^+ + hole \rightarrow Ag^{++}$$
, (1.2)

$$Ag^{+} + Ag^{0} \rightarrow Ag_{2}^{+} \qquad (1.3)$$

The paramagnetic resonance lines of each center are well known. The Aq⁺ ion is not paramagnetic.

The interest of the study of these centers in different matrices rests in the interpretation of the quantitative data, which are sensible to the nature of the **site** where the silver ion is stabilized.

In addition, **two** more kinds of EPR lines were detected: a) boron hole centers²

b) boron electron centers (BEC)¹.

The Ag_2^+ centers were not detected, even with subsequent thermal bleaching.

2. EXPERIMENTAL

The samples were X-irradiated at 77° K during 8 hours, with 45 KV and 35 mA beam using a Phillips unity, type 1008/65, with tungsten target.

The glass samples were given by A. Bishay during hispermanence in the Instituto de **Física** of University of São Paulo. The details of sample production are given in Ref. 3. All the glasses studied contain 30% of BaO and silver enters in the composition with 5% mole. The Al₂O₃ and B₂O₃ content is variable from sample to sample, as indicated in the tables.

The EPR spectra were obtained with X-band reflection homodine spectrometer, using rectangular cavity which operates in the H_{102} mode. The temperature control was done using a Varian unity, model V-4540.

3. TYPICAL SPECTRUM OF BaO; AI,O,; B,O, GLASSES CONTAINING SILVER

In Fig.1 there is a typical spectrum of atomic centers of silver induced by radiation. Care was taken, during the transfer of the sample to the spectrometer, to prevent temperature raising, avoiding the recombination of the more unstable centers.

The assignment of the lines located respectively in 2920 Qe and 3500 Qe to Ag^0 centers needs a careful verification on account of the similarity in position and shape with the EPR lines of the Ag_2^+ center discoved by Mel'nikov *et al.*⁴ in glassy B_2O_3 containing silver impurity.

The Ag⁺⁺ center presents EPR lines with noticeable local anisotropy verified by means of the separate values of $g_{||}$ and g_{\perp} of the splitting factor. To each anisotropic g-factor component corresponds a non--resolved doublet, their positions being in agreement with those studied by Shields⁵ in X-irradiated frozen silver salts. In Fig. 1 the $g_{||}$ doublet corresponds nearly to the values between 3080 Qe and 3100 Qe. The $g_{||}$ doublet is situated in the region of 2800 Qe.



Fig.i - Typical EPR spectrum of borate glass containing 30% BaO and 5% mole of silver, X-irradiated at 77°K.

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In 3.2 k0e a superposition of BR lines appears whose centers are not attributable to silver ions. These centers were identified **res**-pectively as boron hole center and boron electron center.

4. ATOMIC CENTER OF Ag^o

There are two types of silver isotopes with almost thesame relative abundance. Both isotopes, 48.1% of ¹⁰⁷Ag and 51.9% of ¹⁰⁹Ag, have the nuclear spin I = 1/2, giving a doublet of each of them. Both doublets are observable in a lot of materials containing silver, giving four resolved EPR lines. The internal pair is related to ¹⁰⁷Ag isotope and the external to ¹⁰⁹Ag.

The energy levels (see Fig.2) are calculated using the Breit--Rabi formula 6 which, applied to Ag⁰, is

$$W(F,m) = -\Delta W/4 - \mu_{I(\beta_{N})} \beta_{N}^{H_{0}m} \pm (\Delta W/2) \sqrt{1 + 2mx + x^{2}}$$
(4.1)

where

$$\begin{aligned} x &= \left(g_J - \frac{{}^{\mu}I\left(\beta_N\right)}{I} - \frac{\beta_N}{\beta_0}\right) \frac{\beta_0 H_0}{\Delta W}; \\ \Delta W &= A = h\Delta v; \\ {}^{\mu}I\left(\beta_N\right) &= \text{nuclear magnetic moment in nuclear magnetons}; \\ \beta_N &= \text{nuclear magneton }; \\ \beta_0 &= \text{Bohr magneton }; \\ H_0 &= \text{applied magnetic field in Oe}; \\ F &= 1,0; \\ m &= 0 \text{ for } F = 0 \text{ and } m = -1, 0, +1 \text{ for } F = 1; \\ g_J &= -\frac{\mu_J}{J\beta_0}. \end{aligned}$$

The constants $\mu_{\mathcal{I}}(\beta_{\mathcal{N}})$, $\mathcal{G}_{\mathcal{J}}$, β_0 and $\beta_{\mathcal{N}}$ are found in tables of Ref. 7. The transitions corresponding to the observed Ag^0 lines in the spectra are respectively $(F = 1, m = 1) \rightarrow (F = 0, m = 0)$ and $(F = 1, m = 0) \rightarrow (f = 1, m = +1)$.



Fig.2 - Sketch of Ag^0 center energy levels calculated from the formula of Breit-Rabi⁶.

For these transitions, using the formula of Breit-Rabi, we get

$$v = -\Delta v \left(\frac{1}{2} \left(1 + x_1^2\right)^{1/2} + \frac{1}{2} \left(1 - x_1\right) + \frac{2\mu I(\beta_N) \beta_N H_1}{\hbar \Delta v}\right)$$
(4.2)

$$v = -\Delta v \left(\frac{1}{2} \left(1 + \frac{1}{x_2^2}\right)^{1/2} - \frac{1}{2} \left(1 + \frac{1}{x_2}\right) + \frac{2\mu I(\beta_N) \beta_N H_2}{\hbar \Delta v}\right)$$
(4.3)

where \mathbf{v} is the klystron frequency.

The values of $\overline{g_J}$ and $\overline{\Delta v}$ which satisfy simultaneously the equations (4.2) and (4.3) (using the constant $\overline{\mu T}(\overline{\beta N})$ equal to the arithmetic mean between the respective values of nuclear magnetic moment of reach silver isotope) are expressed in Table 1, where the line width of each transition is included.

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Sam	nple	%Al ₂ 0 ₃	*в ₂ 0 ₃	$\delta(\overline{\Delta v})/\overline{\Delta v}_{free}^{\ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ $	Δν . 10 ⁹ Hz	$\overline{g_J}$	Line wid (1,-1)→(0,0)	dth, Ce (1,0)→(1,1)
	E-9	-	70	- 14.6 ± .1	-1.618 ± .022	1.9997 ± .0015	64	52
(a)	E-9	-	70	- 18.7 ± .1	-1.670 ± .022	1.0015 ± .0017	57	52
	E-7	5	65	- 15.7 ± .1	-1.597 ± .022	1.9888 ± .0015	64	48
(Ь)	E-3	20	50	- 19.7 ± .1	-1.522 ± .021	1.9838 ± .0014	64	52
	E-2	20	50	- 18.7 ± .1	-1.536 ± .022	1.9811 ± .0014	68	56
(a.t.)	E-2	20	50	- 22.2 ± .2.	-1.474 ± .029	1.9853 ± .0025	60	48

(a) Repetition of the measure 2 days after, in the same conditions. (b) Sample produced in oxidizing conditions. Very pure $A\ell_2 0_3$ content. (a. t.) Measured at ambient temperature.

The spectrum in Fig. 1 is very similar to the spectrum of Ag_2^+ obtained by Mel'nikov *et al.*⁴, mainly because of the overlapping of the central line of Ag_2^+ triplet with a lot of lines originated *from* boron centers.

In order to eliminate the doubt whether the line located in 2920 Qe and 3500 Qe are associated to Ag_2^+ centers, we observe in detail the central region located in the range of 3080 Qe to 3220 Qe (see Fig. 3). There is a superposition of the spectra of the Ba0 (30%); B_2O_3 (70%) glass containing 5% mole of silver and of the respective base glass. The lines situated in the central part are due to the boron hole center. In the spectrum of glass containing silver the only additional lines aredue to the Ag_{t}^{++} center and the boron electron center assigned as e_{t}^{-} . In base



Fig.3 - Central region viewed in detail of the EPR spectra of thesamples E-8 and E-9 of Ba0(30%); B₂0₃(70%) glasses put together in order to compare X-irradiated sample containing silver with base glass. All the lines present were identified being evidenciated the absence of the third line of Ag⁺₂ center which might be present.

glass the line corresponding to e_t by boron atom were bleached out. It is evident, therefore, the absence of any possible EPR line attributable to Ag_2^{\dagger} center in the central part of the spectrum.

The factors governing the width and the separation of the \mbox{Ag}^0 EPR lines are;

a) concentration of the alkali modifier

b) concentration of silver impurity

The values of the Table 1 are in excellent agreement with the results of Asabghy et $al.^8$, who plotted the hfs constants and the line widths in function of the concentration of alkali oxide in the composition of **bo**-rate glasses containing sodium oxide.

5. ATOMIC CENTER OF Ag++

The Ag^{++} centers are formed by hole trapping of Ag^{+} ions. The interactions of Ag^{++} ions are studied by means of the spin Hamiltonian with axial symmetry^{5,9} in first approximation

$$H = g ||_{0}^{\beta} B_{z} S_{z} + g |_{0}^{\beta} (H_{x} S_{x} + H_{y} S_{y}) + AI_{z} S_{z} + B(I_{x} S_{x} + I_{y} S_{y})$$
(5.1)

where A and B are the hfs constants, g_x = g_y = g_{\perp} and g_z = g_{\parallel} .

Taking 9 as the angle between the local symmetry axis and the direction of applied magnetic field and $g_0 = (g_{||} + 2g_|)/3$, it follows:

$$H_{(\theta=0)} = (g_0 H_0 / g_{\parallel}) - (m_I A / \beta_0 g_{\perp}) ; \qquad (5.2)$$

$$H_{(\theta = \frac{\pi}{2})} = (g_{0}H_{0}/g_{\perp}) - (m_{I}B/\beta_{0}g_{\perp}); \qquad (5.3)$$

where $H_0 = hv/g_0\beta_0$.

Ttie Ag⁺⁺ EPR lines of borate glass containing 30% BaO and 5% mole of silver are indicated in Fig.1, corresponding to non-resolved doublets for eech g₁ and g_{11} . In Fig. 4 the doublet corresponding to g_{11} is



partially resolved. The set of resolved EPR lines of Ag^{++} were obtained by Shields⁵ using frozen D_2O matrix. The same transitions were observed by Yokota et $al.^9$ after thermal bleaching in Li - Al-rnetaphosphate glass containing Pb and Ag. In the last case the resolution of doubletswasnot achieved on account of the multiplicity of types of Ag⁺⁺ sites in the glass structure.

The partial resolution of the doublet of $g_{||}$ in Fig. 4 is very unstable and readily disappears at a temperature of about -100°C during 10 minutes. It was not possible to evaluate the hfs constants A and B of spin Hamiltonian because of the lack of resolution of these doublets. The values of $g_{||}$ and g_{\perp} could evaluated writing the equations (5.2) and (5.3) in the reduced form:

$$H_0 || = \frac{hv}{g_{||}\beta_0}$$
(5.4)

$$H_{0\perp} = \frac{h_{\nu}}{g_{\perp}^{\beta_{0}}}$$
(5.5)

where $H_{0||}$ and $H_{0|}$ correspond to the mean position of the central positions of Ag⁺⁺ doublets.

The values of g and g evaluated for the samples are listed in Table 2.

Sam	ole	%AL203	%B ₂ 0 ₃	g	g⊥
	E-9		70	2.327 ± .017	2.108 ± .013
(a)	E-9		70	2.328 ± .017	2.108 ± .013
	E-7	5	65	2.308 ± .016	2.079 ± .013
(b)	E-3	20	50	2.301 ± .016	2.062 ± .013
	E-2	20	50	2.301 ± .016	2.062 ± .013
(a.t.)	E-2	20	50	2.273 ± .016	2.064 ± .013

The separation of the lines of the $g_{||}$ doublet of Fig. 4 gives an estimate of hfs constant A which is of the order of

$$A/g_{||} \beta_0 = 44 \pm 4 \ 0e$$
 (5.6)

6. BORONELECTRONCENTER (BEC)

The EPR lines related with electrons trapped e_t by boron atoms are indicated in Figures 1, 4 and 5. These centers were first described by Griscom¹ in recent work using EPR study of borate glass with potassium oxide. The observed lines are part of 4 transitions from the interaction between ¹¹B nucleus, more abundant, with the trapped electron. The ¹¹B nucleus has the spin I = 3/2 and the equal separation between adjacent lines gives the hfs constant A.

The hfs constant of BEC is very sensitive to the local dimension of the site where the electron was trapped. In Table 3 is apparent the difference between the interactions in glasses containing about 5% mole of silver and the base glass. We conclude that the presence of silver impurity atoms exerts great influence on the glass structure, whose boron and oxigen chains are destroyed, giving origin to greater rings, where the electron center suffers a smaller interaction.

Sam	ple	%AL203	%В₂0 3	%Ag	< <i>A</i> > (BEC) ^(0e)
	E-9	-	70	5	56 ± 4
(a)	E-9	_	70	5	56 ± 4
	E-7	5	65	5	56 ± 4
(ь)	E-3	20	50	5	50 ± 4
	E-2	20	50	5	60 ± 4
	E-8	-	70		96 ± 4
	E-1	∖ 20	50	-	96 ± 4

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Fig 5 - EPR spectrum of Ba0(30%); $s_2 o_3$ (30%) base a showing three of the furlines of boron electron cente (BEO). The is of BEO for m_T =

= $+\frac{1}{2}$ was obscured by the soron hole center.

Looking at the Table 3, it is apparent that the concentration of Al?_0 $_{\rm 2}$ exerts only a weak influence on the BEC.

The evaluation of the mean value of hfs constant <A> of BEC was done measuring the separation of the lines corresponding to $m_T = -1/2$ and $m_T = +1/2$ in the glasses containing Ag and between the lines corresponding to $m_T = -3/2$ and $m_T = +3/2$ in the base glasses.

The comparatively greater value of <*A*> for thebaseglasses indicates that the boron-oxygen rings have respectively smaller size.

7. DISCUSSION

The relatively high concentrations of alkali oxide (30% of BaO) and Ag (5% mole) favours the formation of an Ag⁰ center with great line width and without isotopic resolution. Asabghy *et al.*⁸ designated this center as Ag_N^0 and showed that the line width is associated to the non--bridging oxygen content in the matrix.

The occurence of the Ag_N^0 center and the increment of *its* concentration is nearly proportional to the silver content in the sample as shown by these authors, who proposed that the Ag_N^0 center results from the association of Ag^0 with a metal atom to form a binary compound.

Since the Ag_2^{\dagger} centers are absent in the samples with relatively high concentration of alkali oxide and silver, it is to be also expected the absence of Ag_2^0 , because in both cases collisions must occur between silver atoms. This would be possible only in favorable conditions for the migration of silver atoms or ions in these materials. So we conclude that in these samples the migration of silver atoms, if any, is very reduced.

The reduced influence of Al₂0₃ concentration on the spin Hamiltonian parameters is due probably to the little ionic radius of aluminium.

The observed Ag^{++} centers gave parameters in good agreement with the values found in literature, related to different matrices.

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The drastic variation of hfs constant of boron electron center showed that the glass structure was affected severely by the **presence** of 5% mole of **silver**, possibly by the cracking of the boron-oxygen rings in the matrix **chain**, opening greater places **in** similar fashion as the effect of the addition of alkali oxide modifier.

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